

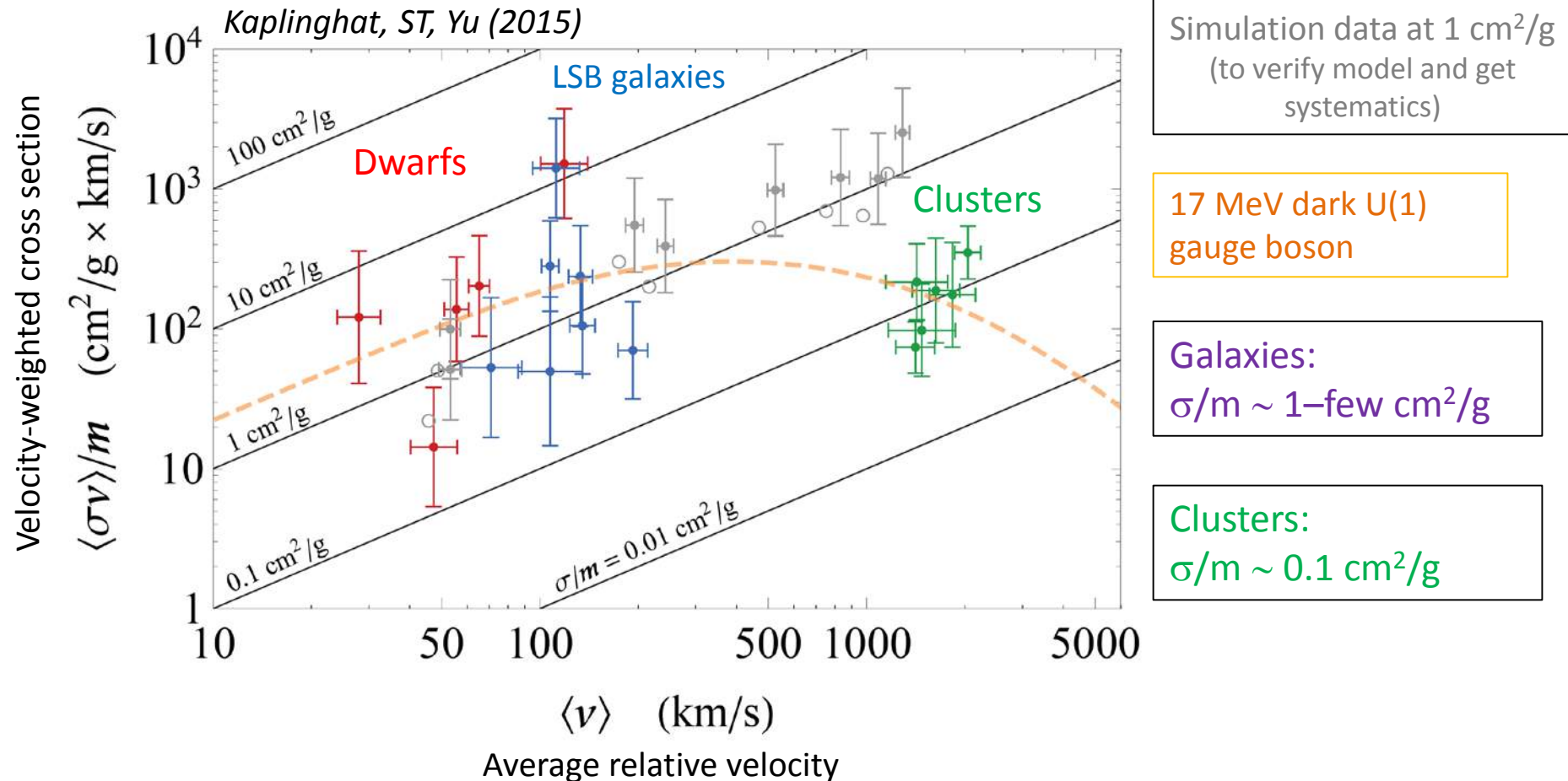
# Complementary constraints on SIDM

Sean Tulin



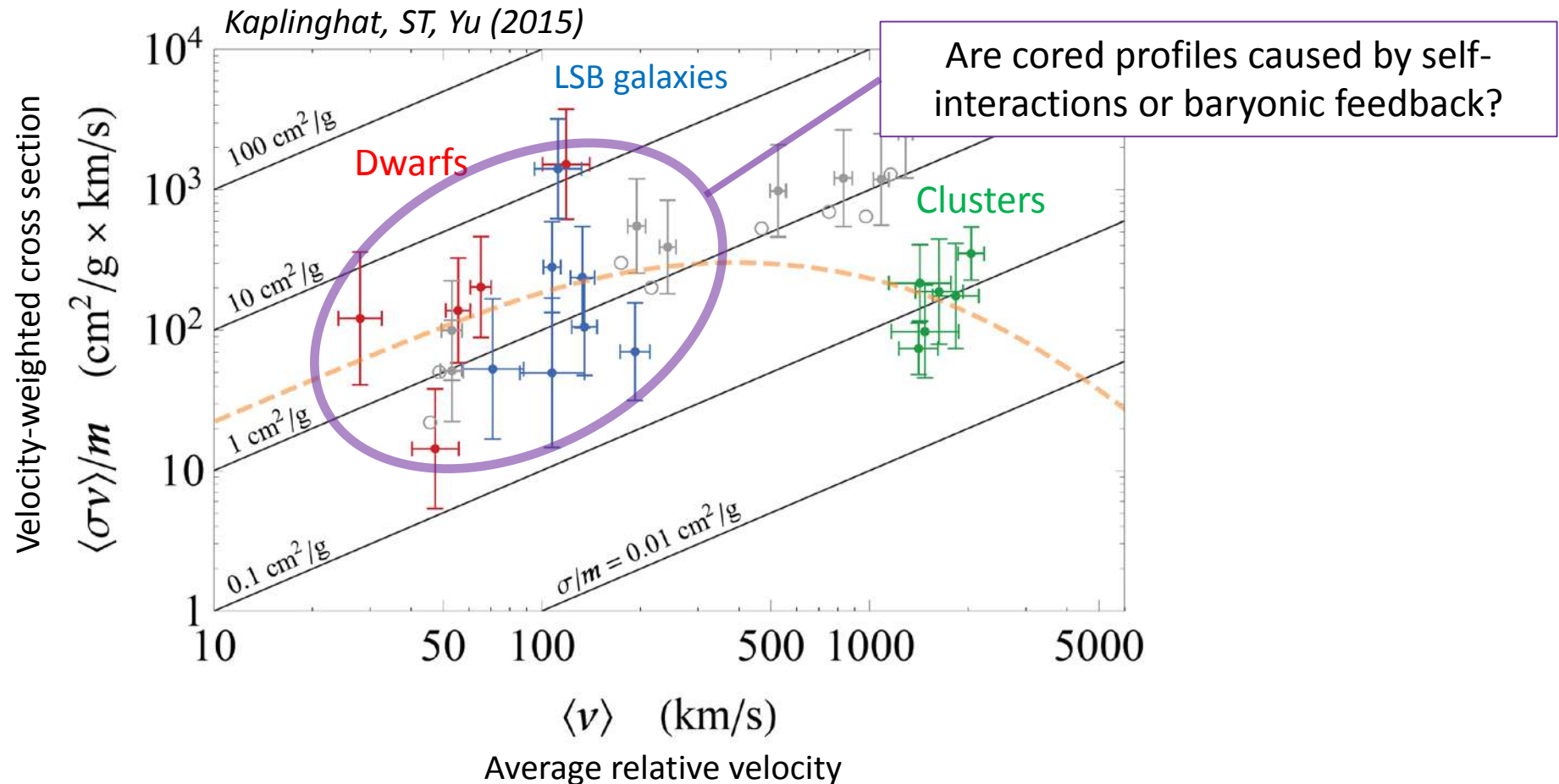
# Velocity-dependent self-interactions

Dark matter self-scattering rate:  $R_{\text{scat}} = \sigma v_{\text{rel}} \rho_{\text{dm}} / m$



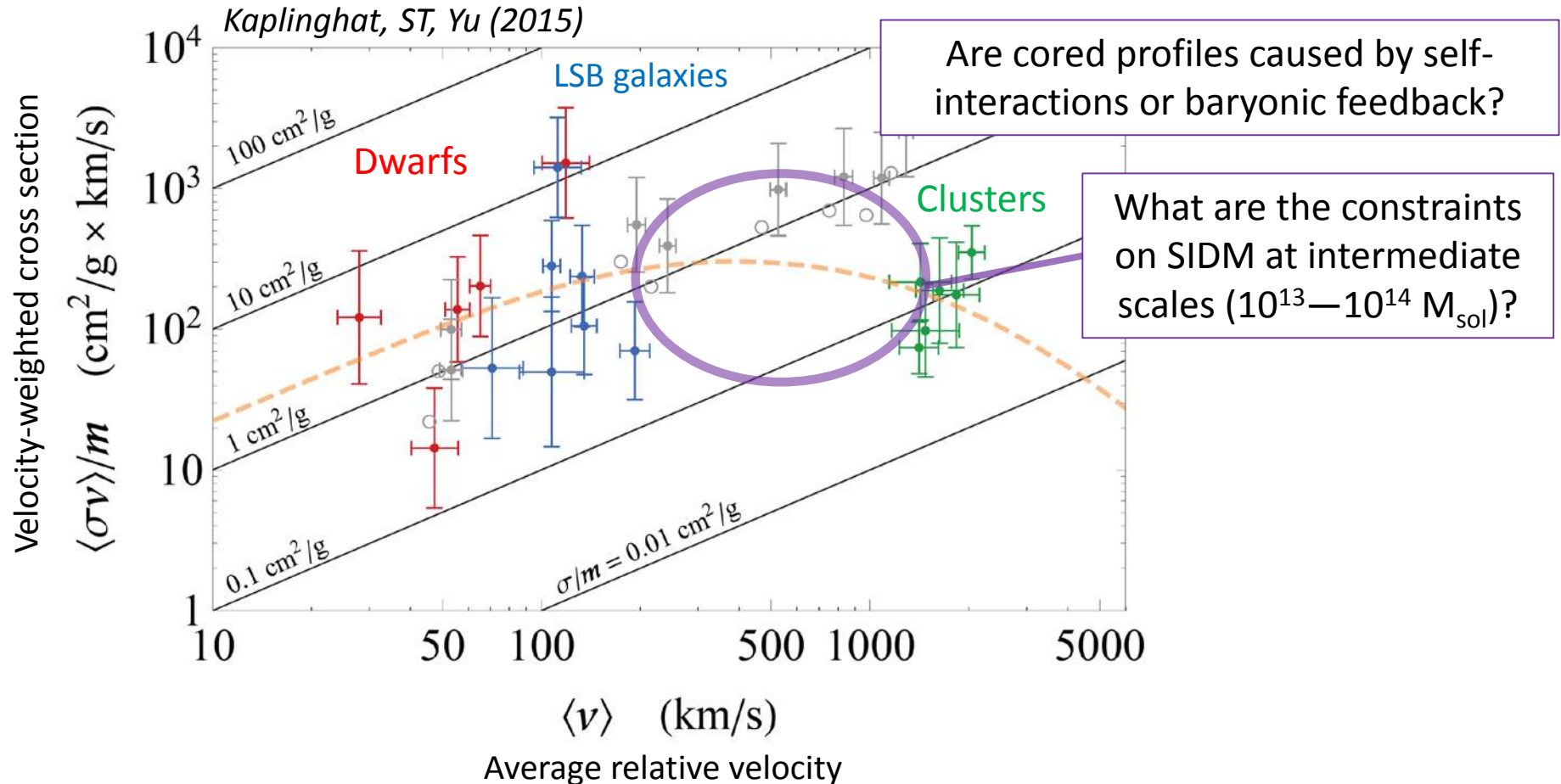
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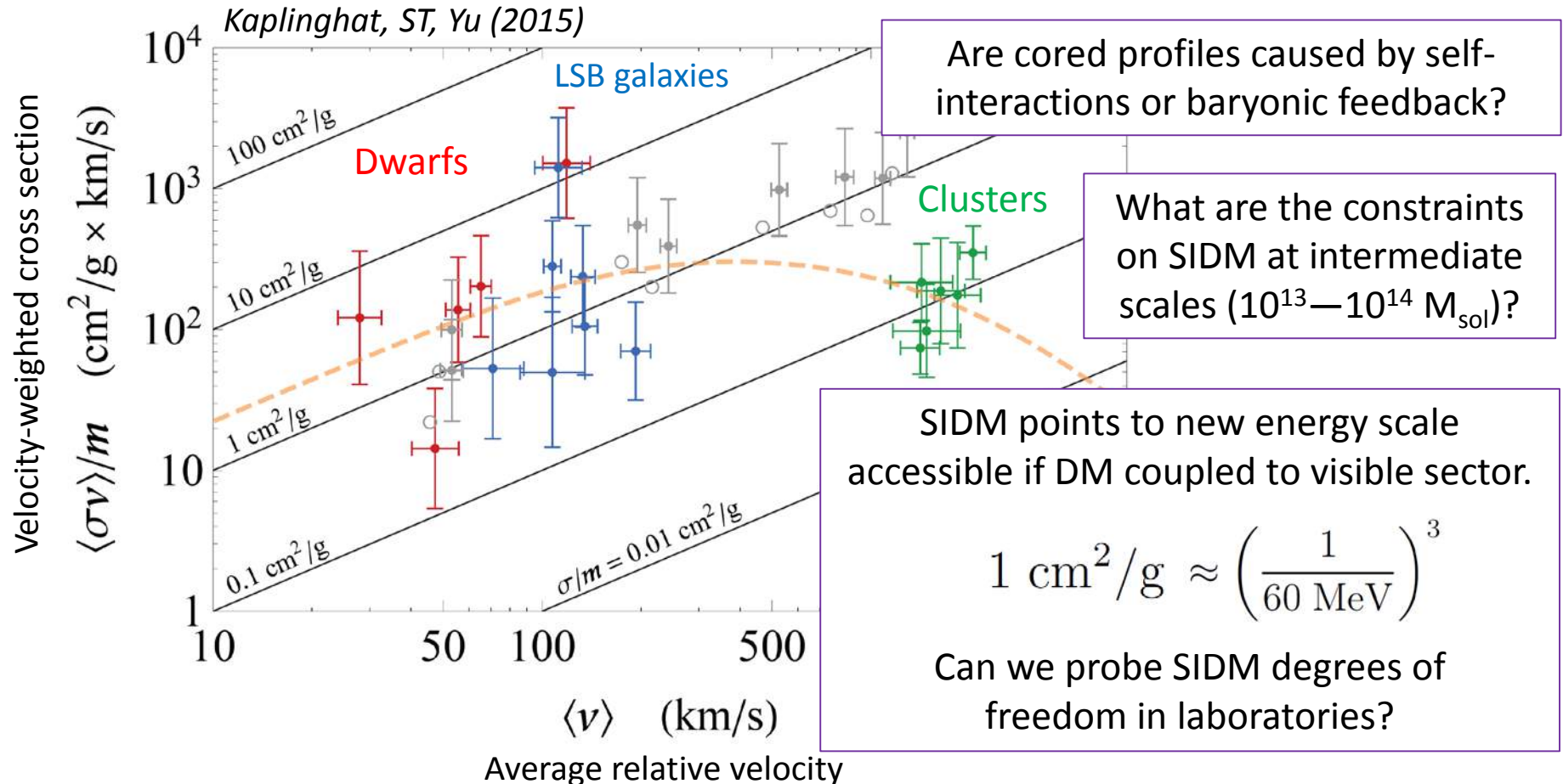
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# Testing bursty feedback

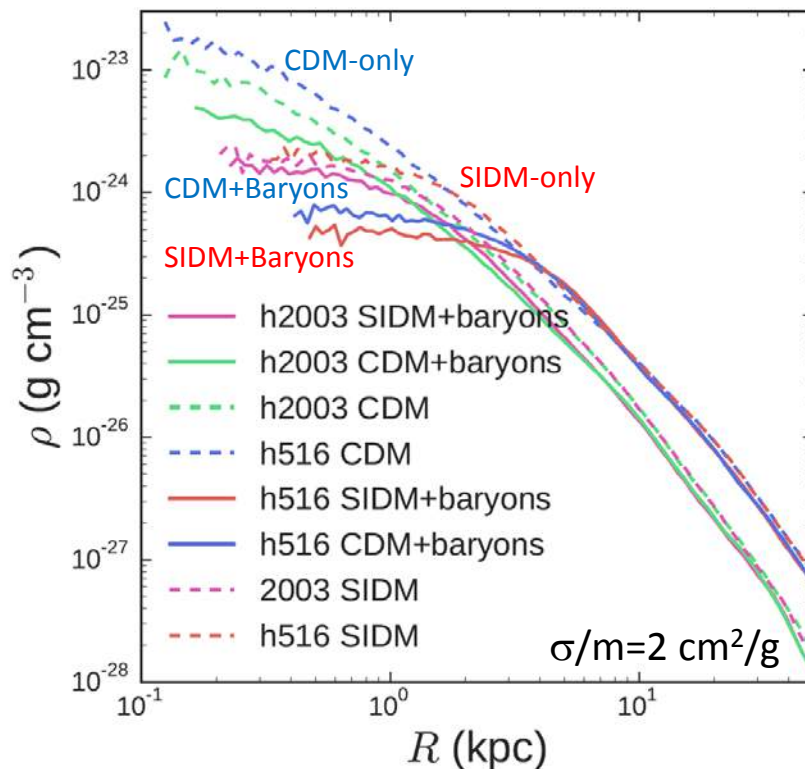
Pressing question: Are all small scale issues solved by feedback?

# Testing bursty feedback

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N-body simulations with SIDM + baryons

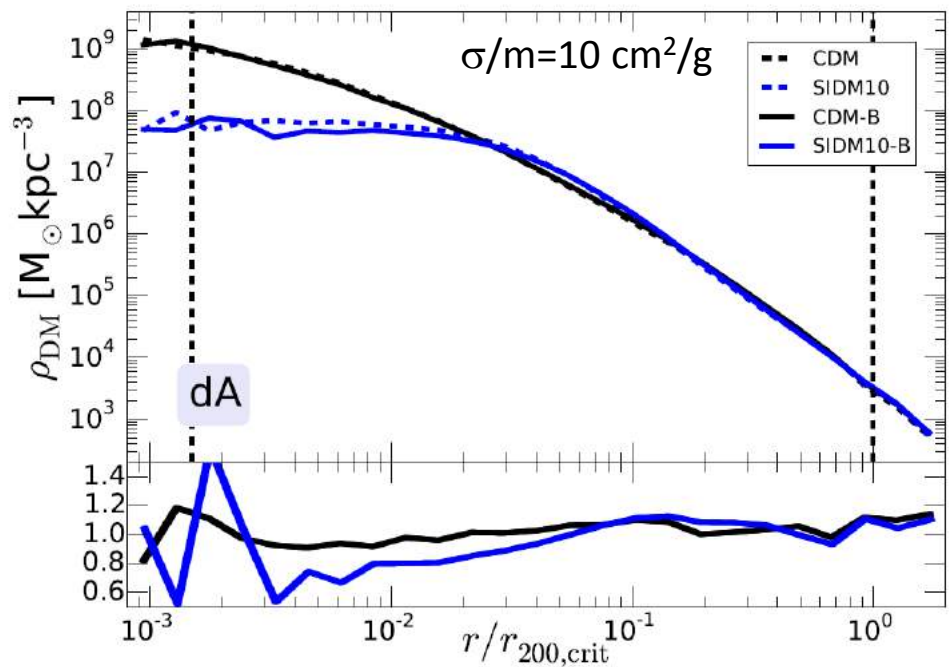
*Bastidas Fry, Pontzen et al. (2015)*



Bursty star formation

(High density threshold for star formation)

*Vogelsberger, Zavala et al. (2014)*



Smooth star formation

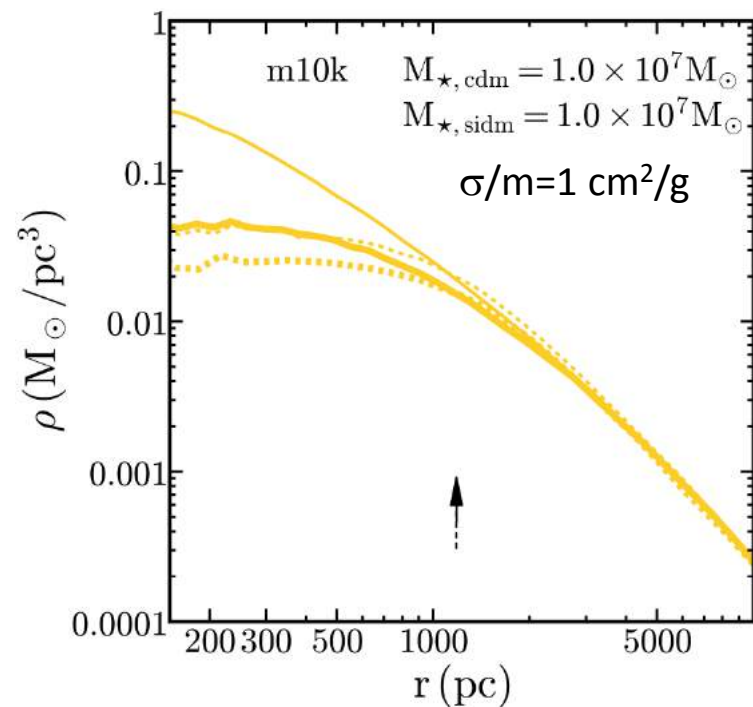
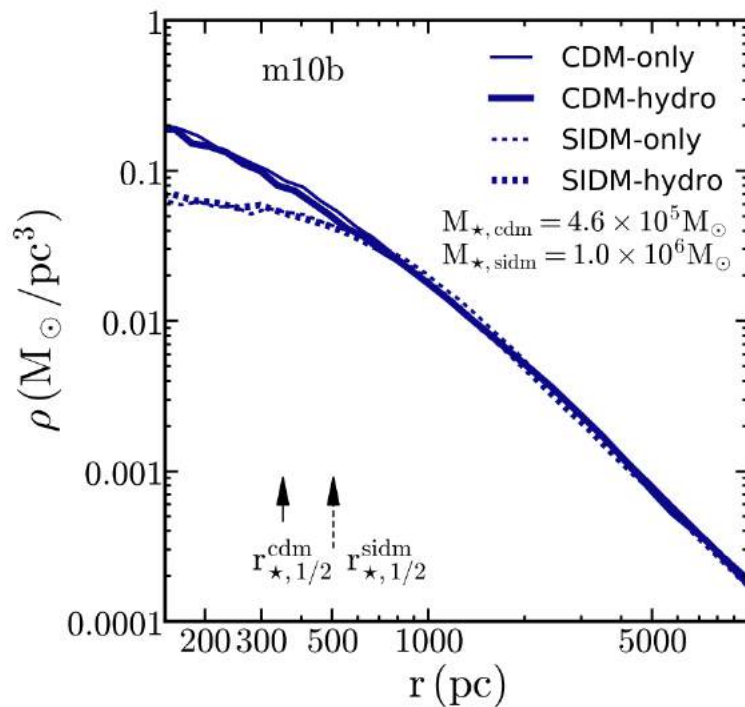
(Low density threshold)

# Testing bursty feedback

Pressing question: Are all small scale issues solved by feedback?

N-body simulations with SIDM + baryons

FIRE simulations: bursty star formation *Robles et al (2017)*





# Resolving timescales

Presence of cored profiles: integrates over 10 Gyr of galaxy formation

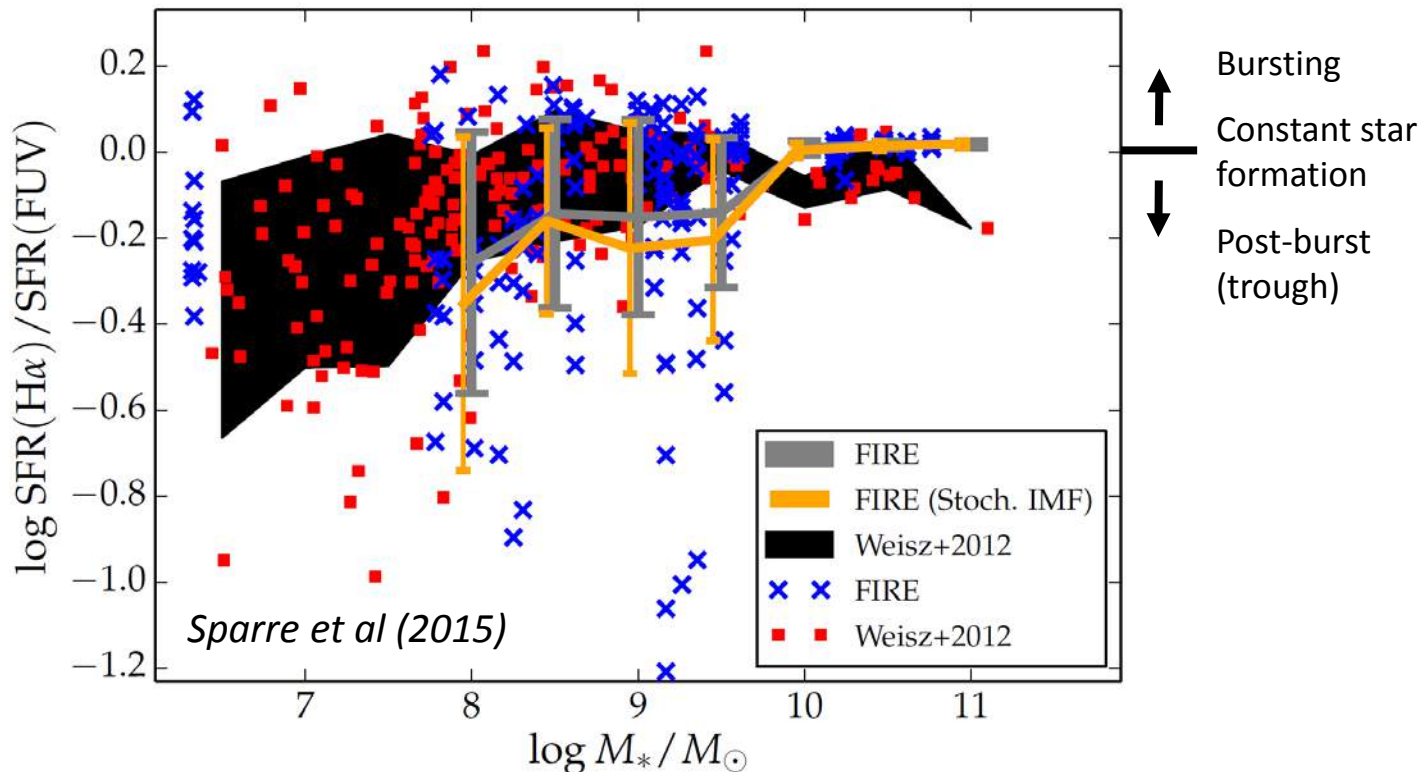
- Nonadiabatic processes acting over  $\sim 10$  Myr (feedback)?
- Adiabatic process over  $\sim 10$  Gyr (SIDM)?

# Resolving timescales

Line ratios as tracer of bursty star formation

$H\alpha$  = traces star formation rate over past 10 Myr

Far-UV = traces star formation rate over past 200 Myr

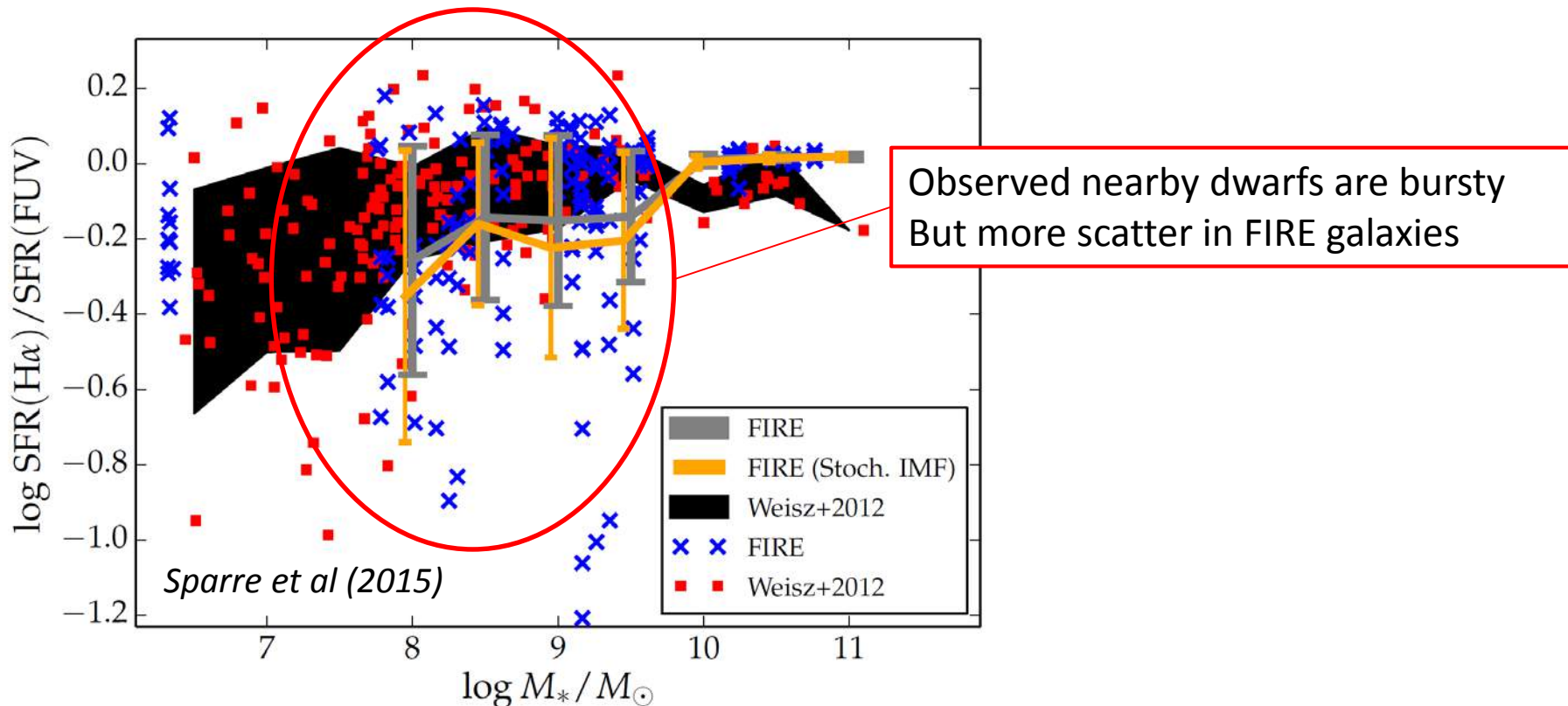


# Resolving timescales

Line ratios as tracer of bursty star formation

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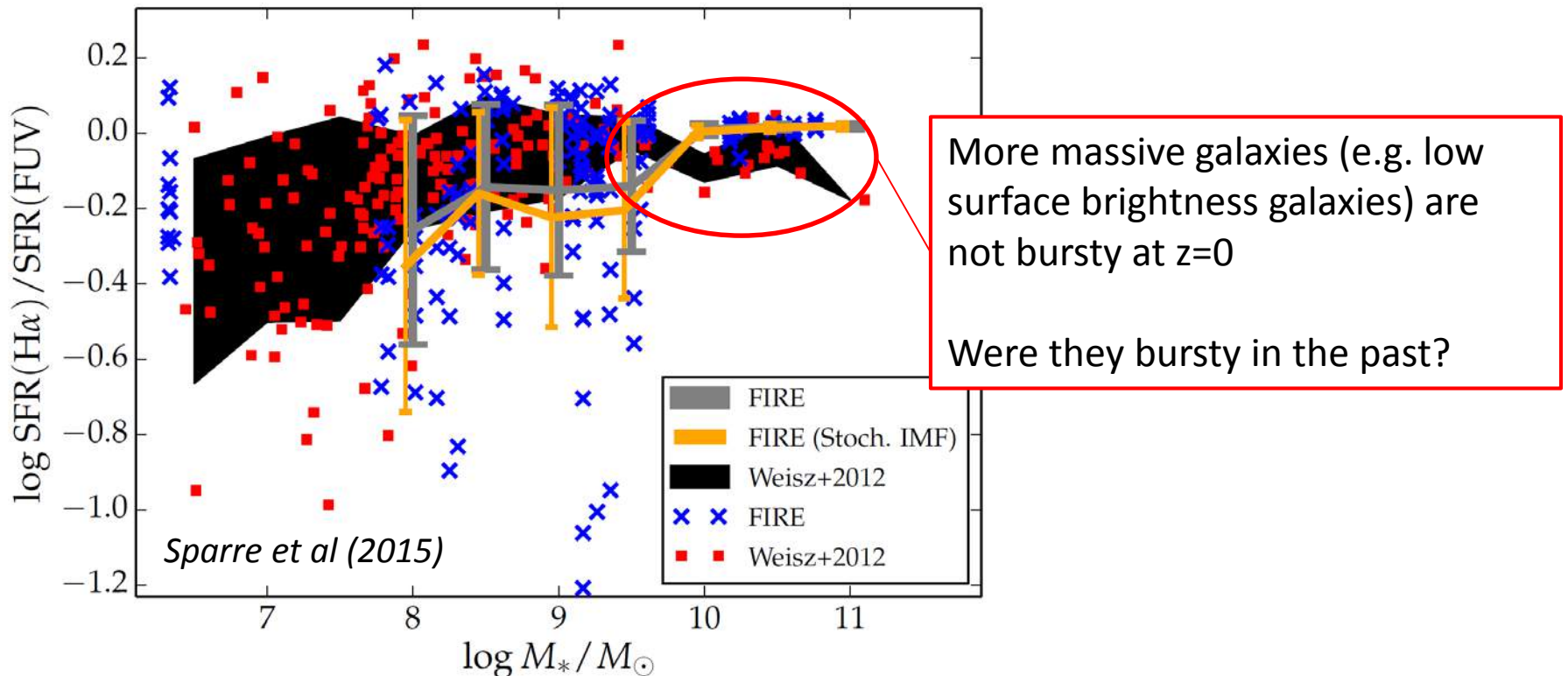


# Resolving timescales

Line ratios as tracer of bursty star formation

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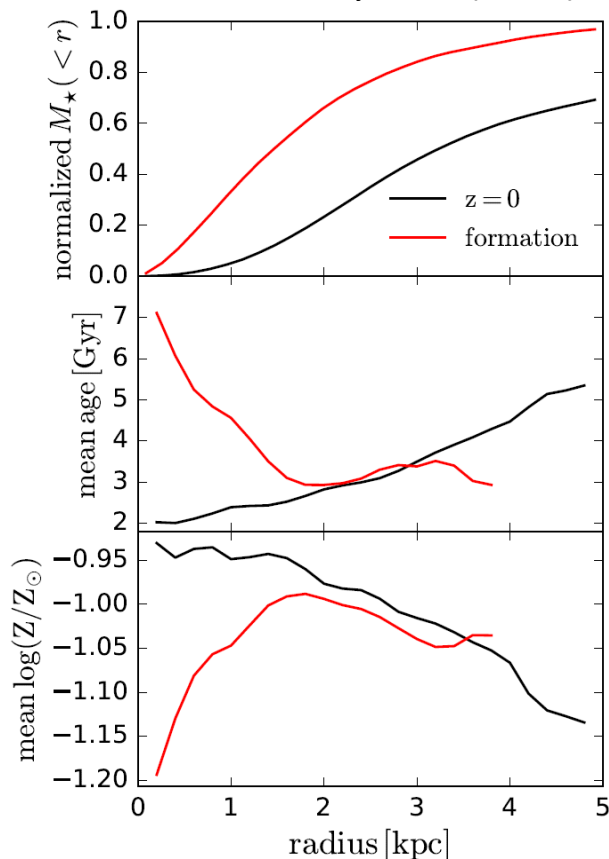
Far-UV = traces star formation rate over past 200 Myr



# Resolving timescales

Star formation history can be imprinted on the radial distribution of stellar populations in nearby galaxies (galactic archeology)

*El-Badry et al (2016)*



Older metal-poor stars form earliest in the center. Bursty feedback causes them to migrate radially outward (along with DM).

Younger metal-rich stars form later throughout the galaxy, but have less time to migrate.

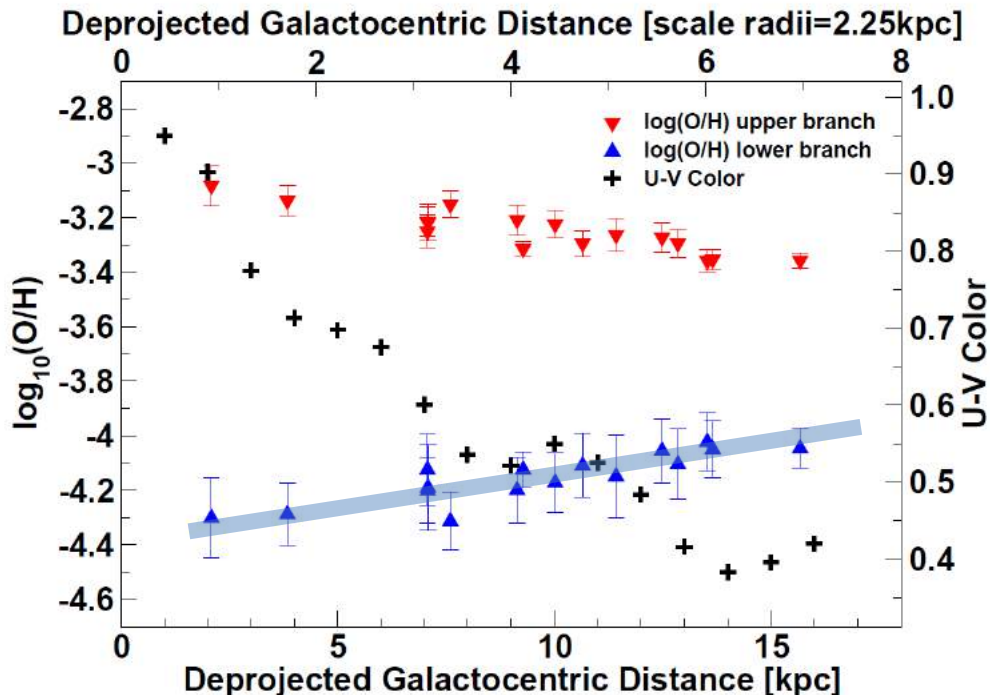
Feedback can **reverse** the age and metallicity gradients in galaxies.

How does this compare to stellar populations in SIDM halos?

# Resolving timescales

Star formation history can be imprinted on the radial distribution of stellar populations in nearby galaxies (galactic archeology)

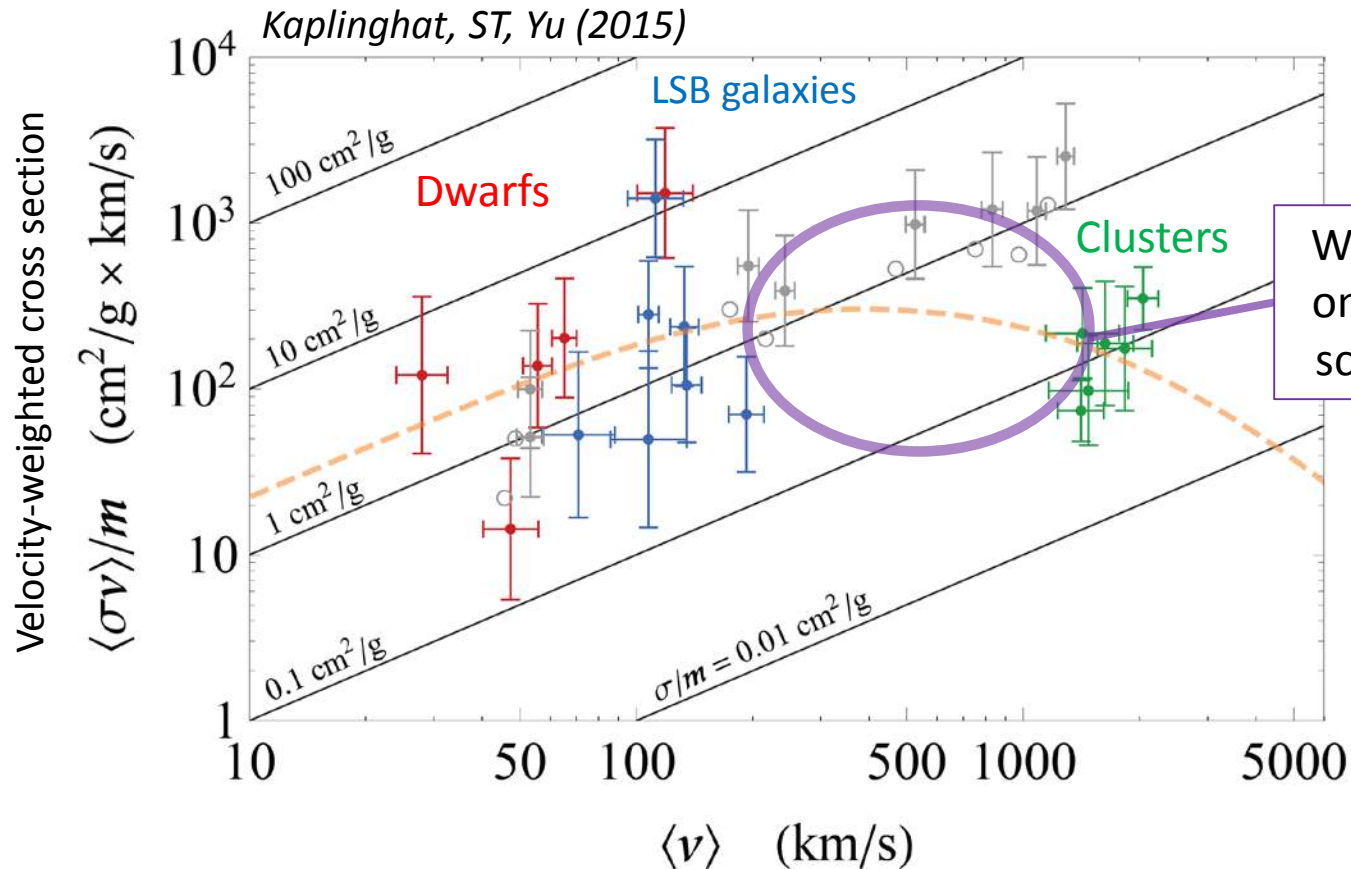
MUltiwavelength observations of the Structure, Chemistry, and Evolution of LSB galaxies (MUSCEL)  
*Young, Kuzio de Naray, Wang (2015)*



Metallicity gradient for LSB UGC 628

Suggests older stars formed earlier in the center, younger stars formed later in the outskirts, with little radial migration.

# SIDM at intermediate scales



What are the constraints on SIDM at intermediate scales ( $10^{13} - 10^{14} M_{\text{sol}}$ )?

Prediction from simple SIDM model: Cross section falls with velocity in this regime.  
Can this be tested?

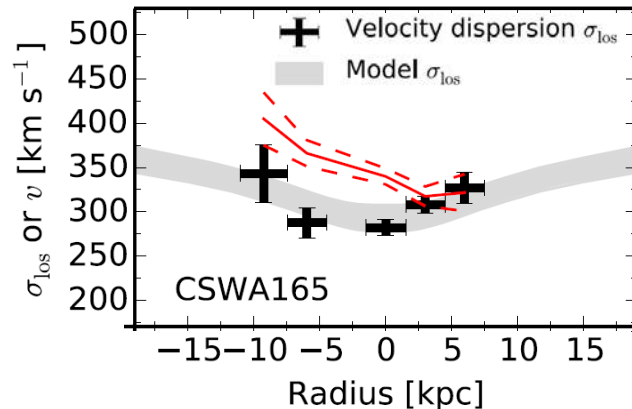
# SIDM at intermediate scales

Mass profile fits of 10 group-scale halos ( $\sim 10^{14} M_{\text{sol}}$ ) with central elliptical galaxies (BGGs)

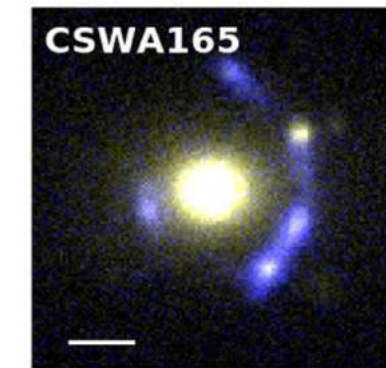
*Newman, Ellis, Treu (2015)*

Data include:

Stellar kinematics in BGG

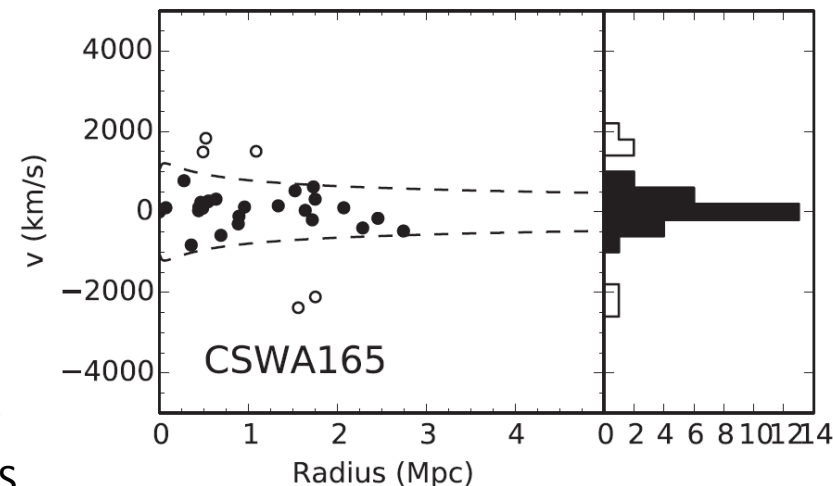


Strong lensing



CASSOWARY survey of lensing systems in SDSS

Kinematics of member galaxies

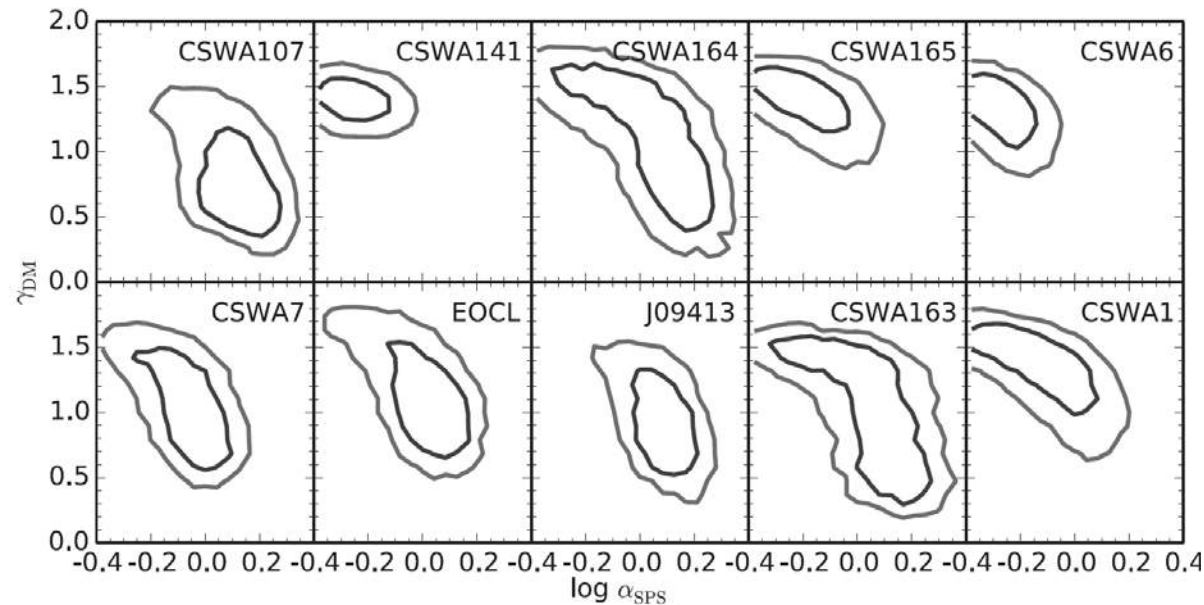


$\sigma$ (km s <sup>-1</sup> )	$\log M_{200}/M_{\odot}$
$362 \pm 52$	$13.78 \pm 0.21$



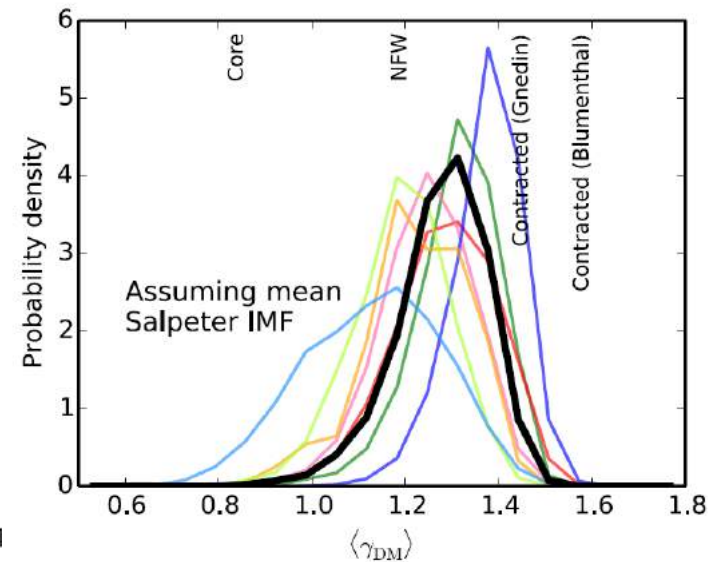
# DM profiles in group-scale halos

Average inner DM slope within inner  $\sim 30$  kpc ( $\gamma_{\text{DM}} \sim 1$  for NFW)



Mass-to-light ratio relative to stellar model  $\alpha_{\text{SPS}} = \Upsilon_V / \Upsilon_V^{\text{SPS}}$

Joint inner DM slope for  $\alpha_{\text{SPS}} = 1$



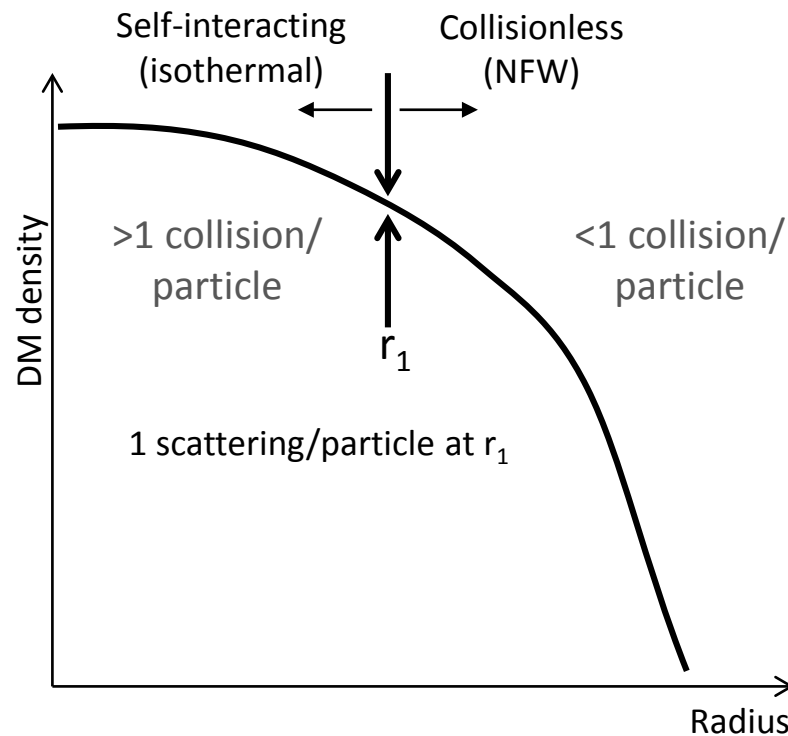
*Newman, Ellis, Treu (2015)*

Halos are more consistent with NFW profiles than cores  
What is the implication for self-interactions?

# Jeans method for SIDM halo profiles

*Rocha et al (2012); Kaplinghat et al (2013); Kaplinghat, ST, Yu (2015)*

Observation-driven approach for deriving density profiles without N-body simulations



Match profiles at  $r_1$ :

$$\rho_{\text{dm}}(r) = \begin{cases} \rho_{\text{iso}}(r), & r < r_1 \\ \rho_{\text{NFW}}(r), & r > r_1 \end{cases}$$

Solve rate equation  
to get cross section:

$$\text{rate} \times \text{time} \approx \frac{\langle \sigma v \rangle}{m} \rho(r_1) t_{\text{age}} \approx 1$$

$$\frac{\langle \sigma v_{\text{rel}} \rangle}{m} = \frac{1}{\rho_{\text{dm}}(r_1) t_{\text{age}}}$$

# Constraints on SIDM cross section

*Sophia Nasr, Laura Sagunski, ST (in progress)*

Scan over halo parameters  $M_{200}$ ,  $c$ ,  $\Upsilon_*$ ,  $r_1$  fitting to data (MCMC).

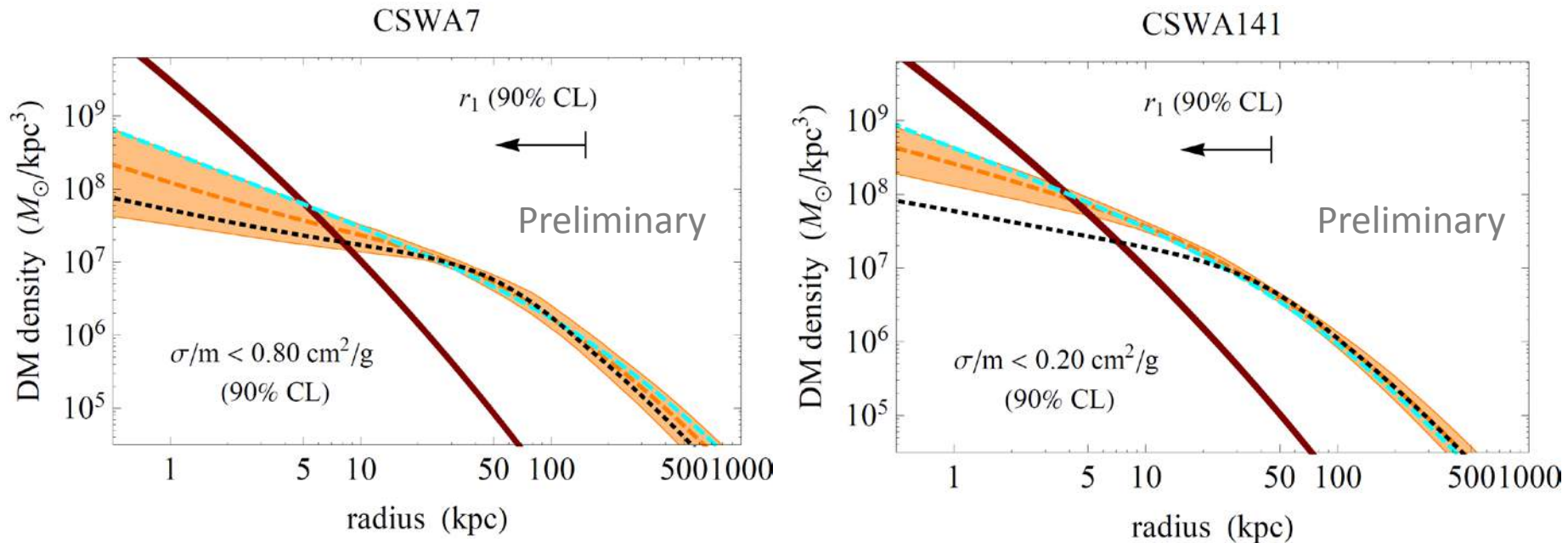
Impose priors (perhaps overly-restrictive):

$(M_{200}, c)$  satisfy mass-concentration relation within scatter

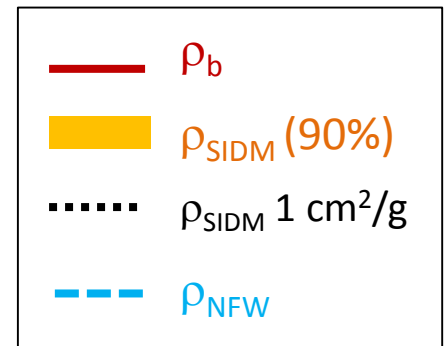
Assume  $\Upsilon_*$  fixed assuming common IMF for all systems within 0.1 dex

# Constraints on SIDM cross section

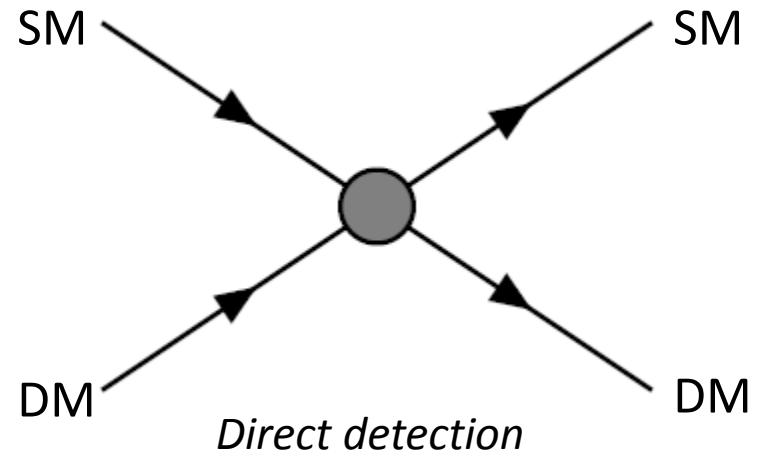
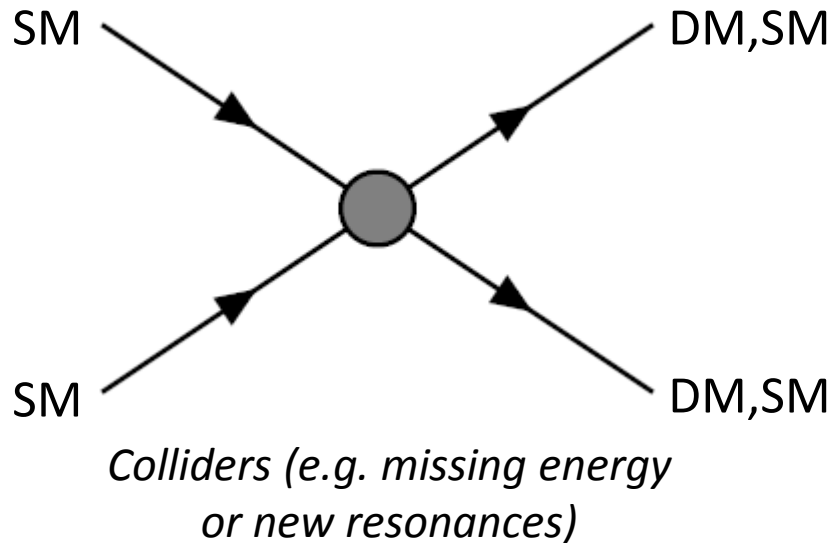
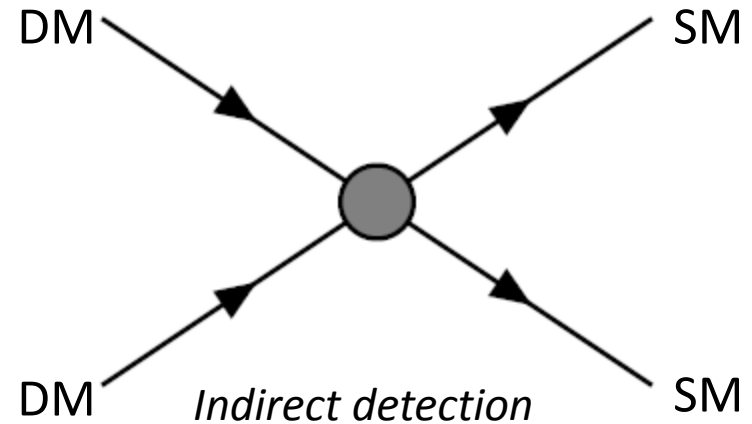
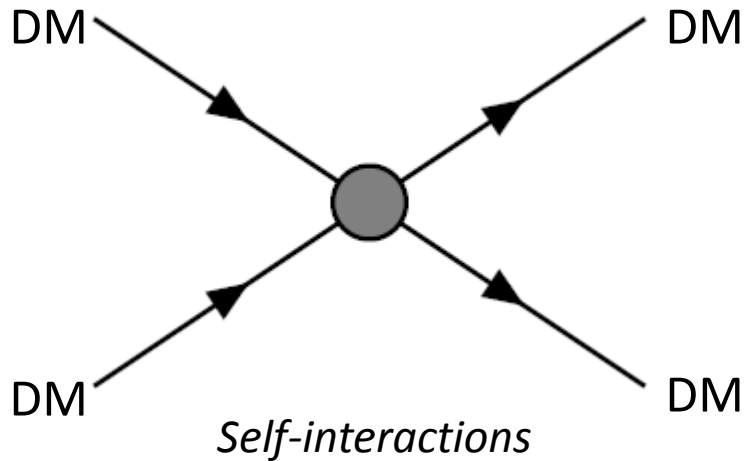
*Sophia Nasr, Laura Sagunski, ST (in progress)*



No conclusions yet, but in principle  
important constraints at  $v \sim 1000 \text{ km/s}$

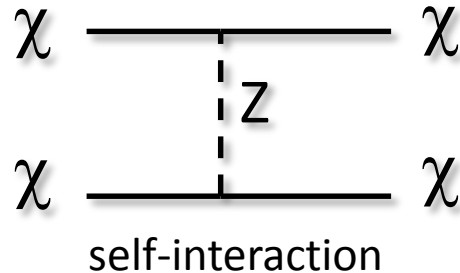


# Complementary probes of SIDM



# Particle physics of self-interactions

WIMPs have self-interactions (weak interaction)



$\chi$  = dark matter (e.g. SUSY particle)

Z boson = mediator particle

Cross section:

$$\sigma \sim \frac{g^4 m_\chi^2}{m_Z^4} \sim 10^{-36} \text{ cm}^2$$

Mass:

$$m_\chi \sim m_Z \sim 100 \text{ GeV}$$

WIMP self-interaction cross section is way too small

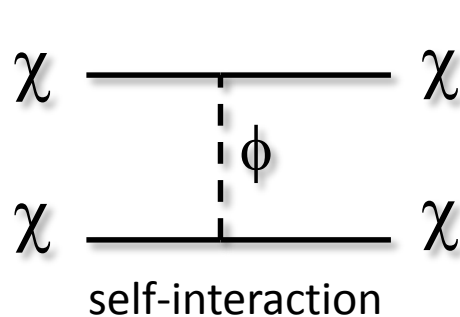
$$\sigma/m_\chi \sim 10^{-14} \text{ cm}^2/\text{g}$$

# Particle physics of self-interactions

Large cross section required  $\sigma/m_\chi \sim 1 \text{ cm}^2/\text{g}$

MSIDM (Minimal SIDM) model: DM + light mediator  $\phi$

*Feng, Kaplinghat, Yu (2009); Buckley & Fox (2009); Loeb & Weiner (2011); ST, Yu, Zurek (2012+13)*



Cross section:  $\sigma \sim \frac{g^4 m_\chi^2}{m_\phi^4}$

Mediator mass below than weak scale

$$m_\phi \sim 1 - 100 \text{ MeV}$$

Velocity-dependence controlled by mediator mass  $m_\phi$

Hard-sphere scattering  
Constant cross section

$$m_\phi \gg m_\chi v_{\text{rel}}$$

Rutherford-like scattering  
Cross section falls with  $1/v_{\text{rel}}^4$

$$m_\phi \ll m_\chi v_{\text{rel}}$$

# What type of models are viable?

How are they testable **beyond** the usual SIDM observables?

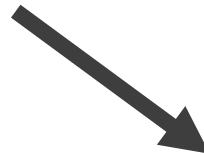
- Light mediator models
- Strongly-interacting DM
  - QCD-like theories
  - Dark hadrons or dark nuclei
- Massless mediator models
  - Dark atoms
  - DM with dark radiation



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Must be largely decoupled  
from visible sector to avoid  
too-large  $N_{\text{eff}}$

*Cosmological probes*  
*See talk by Cyr-Racine*

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Compelling possibilities for  
SIDM to interact with visible  
sector

*Canonical WIMP searches  
(focus of this talk)*

Or may be decoupled  
*Cosmological probes*


Must be largely decoupled  
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- Light mediator models
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May be coupled to visible sector (not as compelling)  
*Canonical WIMP searches*

Or may be decoupled  
*Cosmological probes*

# Strongly-interacting dark matter

Dark sector has non-abelian gauge symmetry

Dark matter is composite made of dark quarks

Lightest baryon  $B_{dark}$  is stable DM due to accidental symmetry (like proton)

May be dark baryon asymmetry ( $\Omega_b \sim 0.2 \Omega_{DM}$ )

# What is the DM baryon $B_{dark}$ ?

Many unknowns:

- Gauge group:  $SU(N)$ ,  $SO(N)$ ,  $Sp(N)$ , ...
- Number of colors  $N_c$
- Number of flavors  $N_F$
- Representations of dark quarks
- Masses of dark quarks  
Chiral or heavy flavor limit?
- Confinement scale  $\Lambda_{DM}$
- Couplings to SM (assume small)

# What guidance do we have?

- Stability

DM baryons are stable

Dark glueballs can also be cosmologically long-lived

- Minimality

- Relic density

Assume  $2 \rightarrow 2$  annihilation  $m_{B_{\text{dark}}} \lesssim 100 \text{ TeV}$

- Astrophysical small scale structure (SIDM)

# Self-interactions for composite DM

Low-energy cross section  $\sigma = 4\pi a^2$

$a$  = scattering length

$m$  = DM mass

$$\sigma/m \sim 3 \text{ cm}^2/\text{g} \times \left( \frac{\Lambda_{\text{DM}}}{m} \right) \left( \frac{\Lambda_{\text{DM}}}{a^{-1}} \right)^2 \left( \frac{100 \text{ MeV}}{\Lambda_{\text{DM}}} \right)^3$$

If dimensionful parameters set by confinement scale

i.e.  $m, a^{-1} \sim \Lambda_{\text{DM}}$ , then  $\Lambda_{\text{DM}} \approx \mathcal{O}(0.2 - 0.6) \times \Lambda_{\text{QCD}}$

gives right  $\sigma/m$  on dwarf scales

# Self-interactions for composite DM

Require velocity-dependent cross section

Naïve condition:

Scattering cross section transitions to being velocity dependent when  $mv_{\text{rel}} \gtrsim a^{-1}$

Since  $v \ll 1$  (DM non-relativistic), require  $ma \gg 1$

More specifically: Want  $\sigma/m$  to transition at  $v/c \sim 10^{-3}$ .

Expect  $ma \sim 10^3$  may provide good fit to galaxies and clusters.



# Self-interactions for composite DM

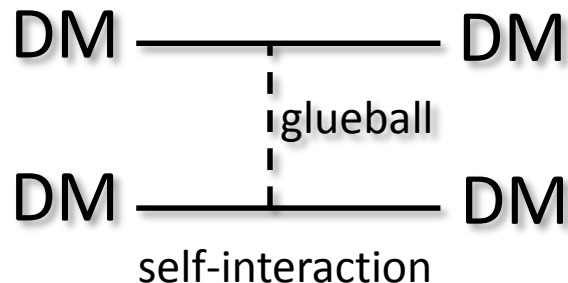
Either  $m$  or  $a$  (or both) must be much larger than naïve scaling with  $\Lambda_{\text{DM}}$ .

- Large mass due to heavy constituents ( $m \gg \Lambda_{\text{DM}}$ ):

Heavy flavor dark quarks

Example:  $\text{SU}(N)$  + heavy adjoint fermion (gluino) *Boddy et al. (2014)*

DM = glueballino (one heavy gluino + gluons)



$$m_{\text{gluino}} \sim 1 \text{ TeV}$$

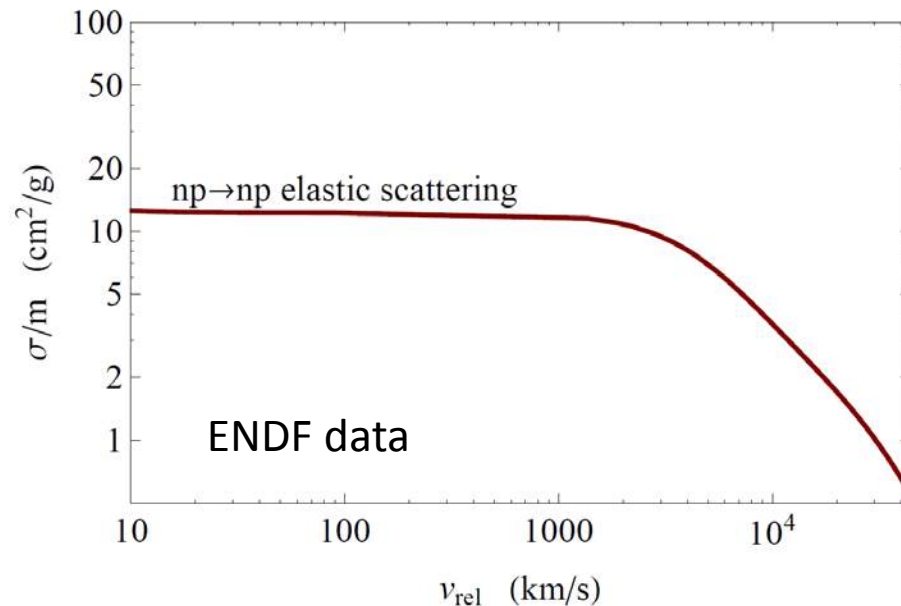
$$m_{\text{glueball}} \sim \Lambda_{\text{DM}} \sim 10 \text{ MeV}$$

# Self-interactions for composite DM

- Large scattering length ( $a \gg \Lambda_{\text{DM}}^{-1}$ ):

Example: nucleon-like DM *Cline et al. (2013)*

Elastic n-p scattering is enhanced due to weakly-bound deuteron  $a^{-1} \approx 15 \text{ MeV} \ll \Lambda_{\text{QCD}}$



Note:

Fall off at  $v \sim c/(ma) \sim 5000 \text{ km/s}$   
Clusters require 0.1 at 1500  $\text{km/s}$

E.g. make  $m_{\text{DM}} \sim 10m_p$  with fixed  
scattering length  $a$

If self-interactions solve astrophysical small scale structure anomalies...

Preferred cross sections

$$\sigma/m \approx \begin{cases} 2 \text{ cm}^2/\text{g} & v \sim 50 \text{ km/s (dwarf galaxies)} \\ 0.1 \text{ cm}^2/\text{g} & v \sim 1500 \text{ km/s (clusters)} \end{cases}$$

Large cross section requires  $\Lambda_{\text{DM}} < 100 \text{ MeV}$

Velocity dependence:

Heavy flavors

Large scattering length

# A minimal theory of DM baryons

*Anthony Francis, R. Jamie Hudspith, Randy Lewis, ST (work in progress)*

Dark QCD-like theory with  $N_c=2$  &  $N_F = 1$   
Fundamental fermion that is SM singlet

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}^a F^{a\mu\nu} + \bar{Q}(i\gamma^\mu D_\mu - m)Q$$

Q = dark quark

F = dark gluon field strength

Possible coupling to visible sector via higher dimensional operators

# What are the lightest states?

Normal QCD ( $N_c=3$ ):

Accidental  $U(1)_B$  baryon number symmetry

Quarks have  $B = 1/3$ , antiquarks have  $B = -1/3$

Dark QCD: weird feature of  $N_c=2$

Fundamental representation of  $SU(2)$  is pseudoreal

Put quarks and antiquarks in unified multiplet

Accidental  $SU(2)_B$  symmetry: *baryon number becomes isospin*

$$\begin{pmatrix} Q \\ \bar{Q} \end{pmatrix}$$

Baryons are isospin multiplets

Lightest baryon is spin-1 (vector) iso-triplet

$$\begin{pmatrix} \rho^+ \\ \rho^0 \\ \rho^- \end{pmatrix} = \begin{pmatrix} QQ \\ \frac{1}{\sqrt{2}}(Q\bar{Q} - \bar{Q}Q) \\ \bar{Q}\bar{Q} \end{pmatrix}$$

# What are the lightest states?

Normal QCD ( $N_f=3$ ):

8 pseudo-Goldstone bosons + massive  $\eta'$  from broken  $U(1)_A$

Dark QCD ( $N_f=1$ ):

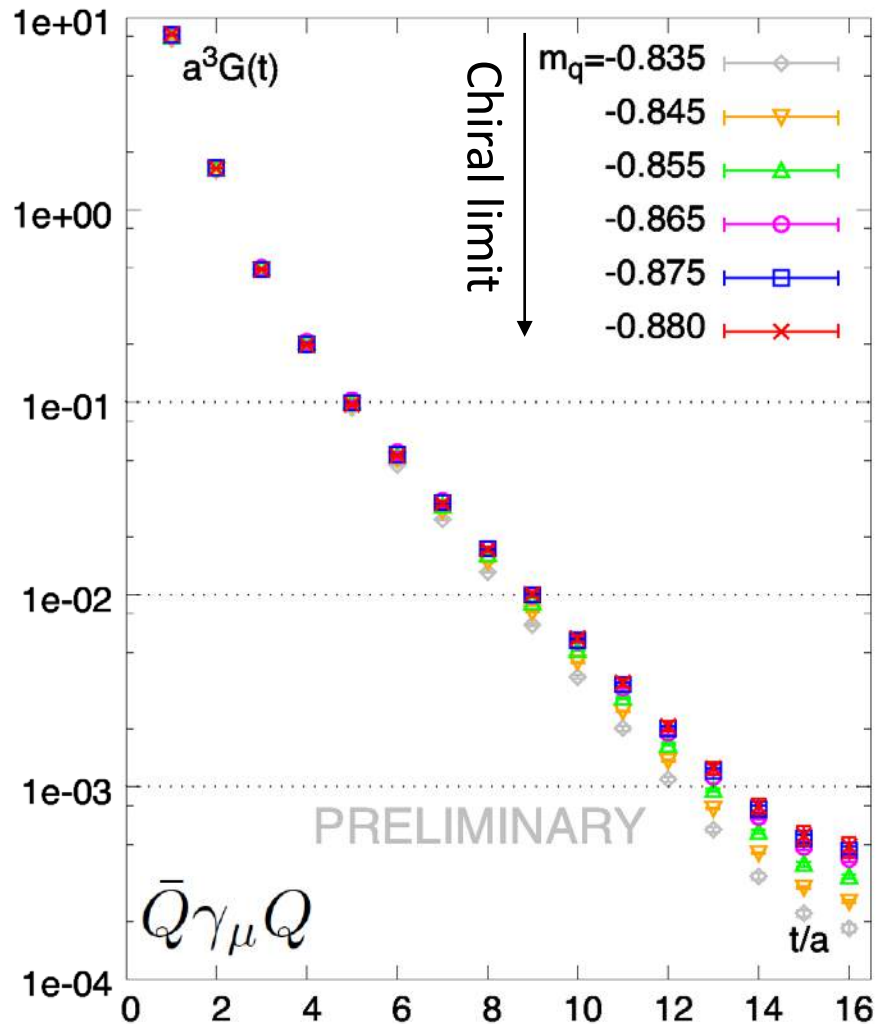
Zero pseudo-Goldstone bosons. Massive  $\eta'$  from broken  $U(1)_A$

Lightest dark states:  $\rho = \begin{pmatrix} \rho^+ \\ \rho^0 \\ \rho^- \end{pmatrix}$   $\eta'$   
vector pseudoscalar

Which state is lighter? Normal QCD:  $m_\rho < m_{\eta'}$

Dark QCD: DM annihilation  $\rho\rho \rightarrow \eta'\eta'$  for  $m_\rho > m_{\eta'}$

# Dark matter on the lattice



Correlator:

Sum over propagators with same quantum numbers as source

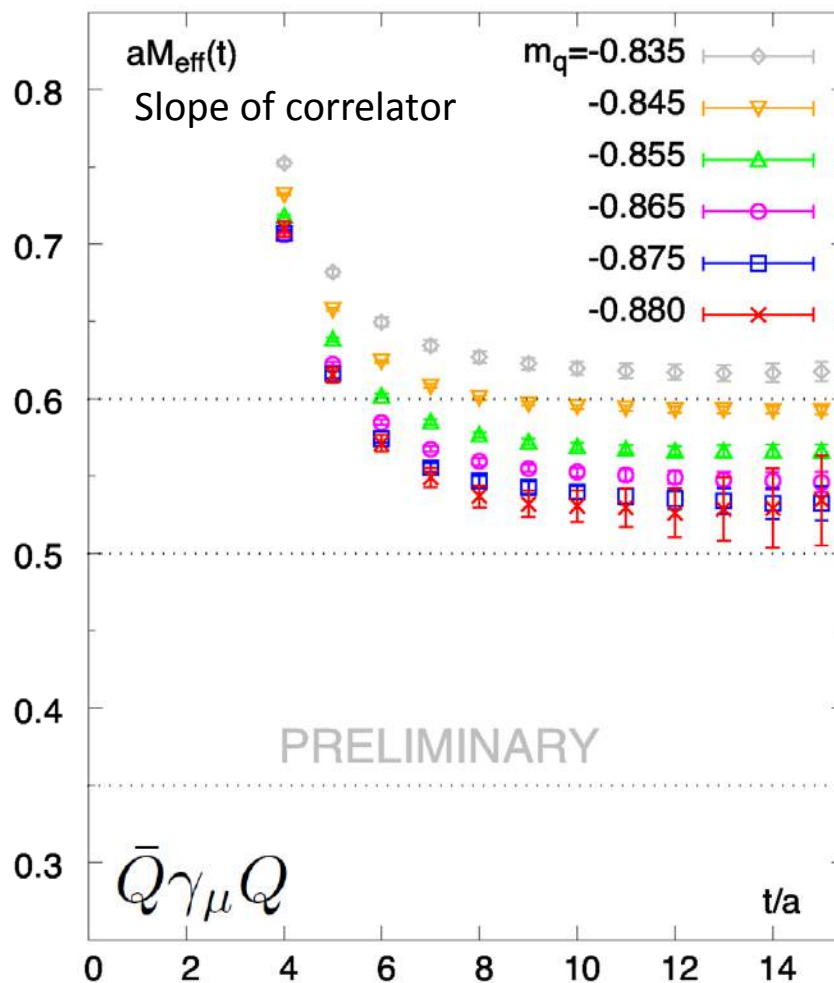
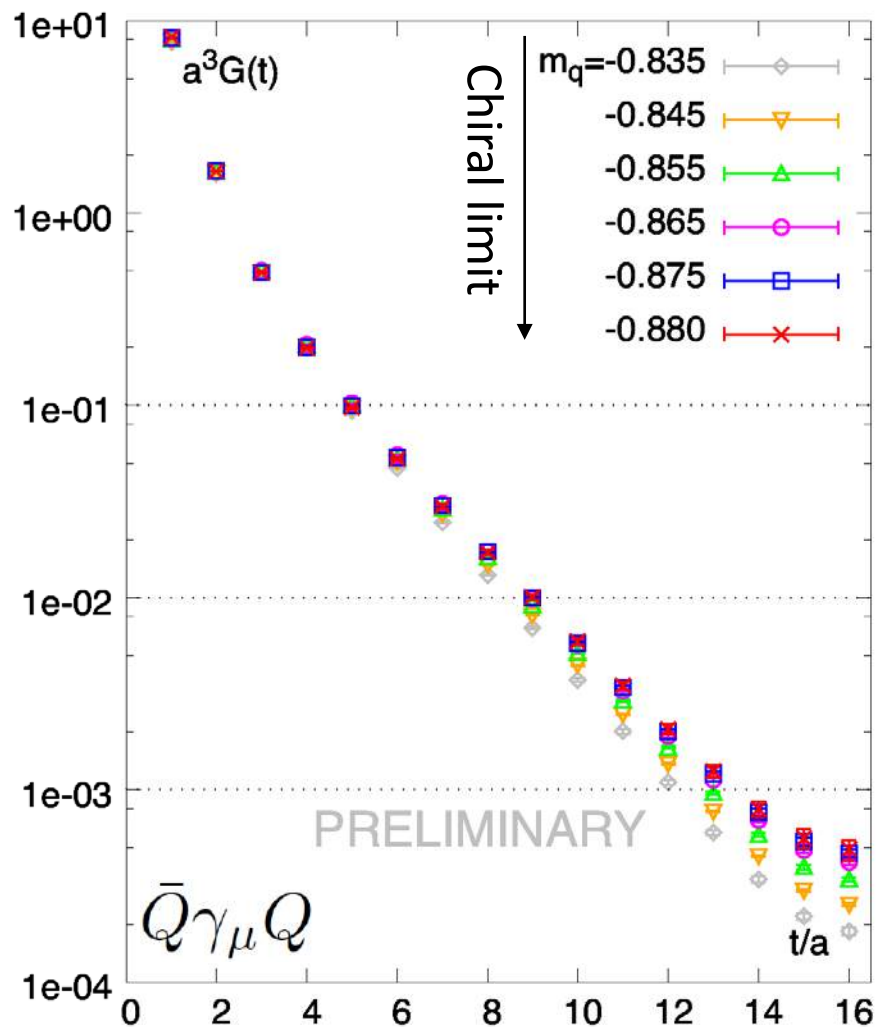
$$G(t) \sim \sum_{\text{states } i} e^{-m_i t} \rightarrow e^{-m_\rho t}$$

Lowest lying state decays slowest (imaginary time)

Log slope at large time gives mass

Anthony Francis, R. Jamie Hudspith, Randy Lewis, ST (work in progress)

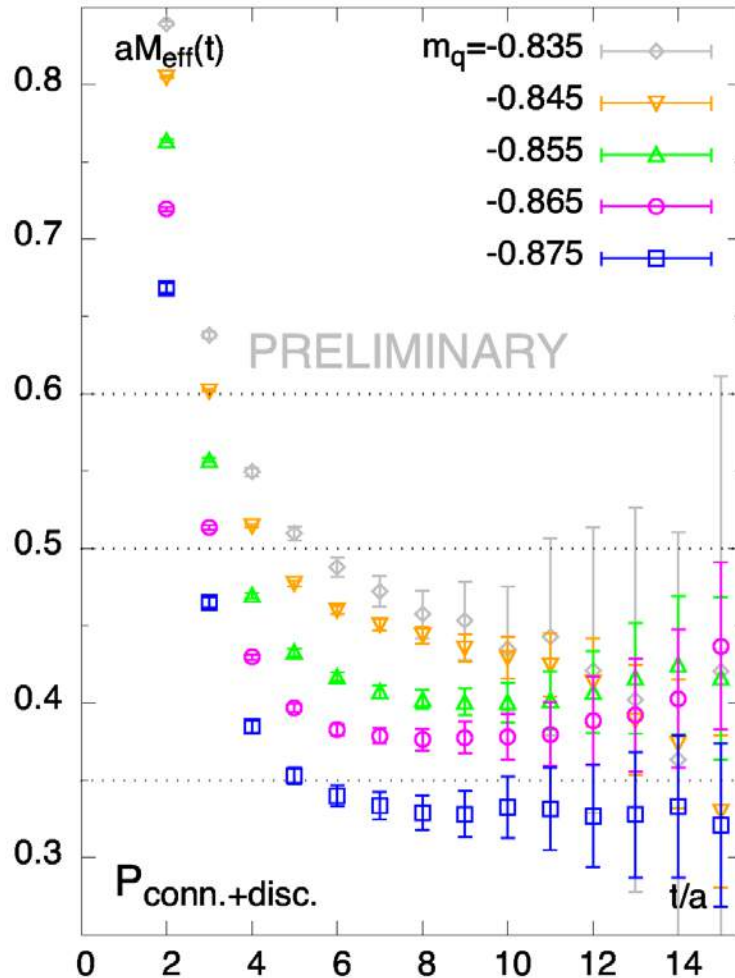
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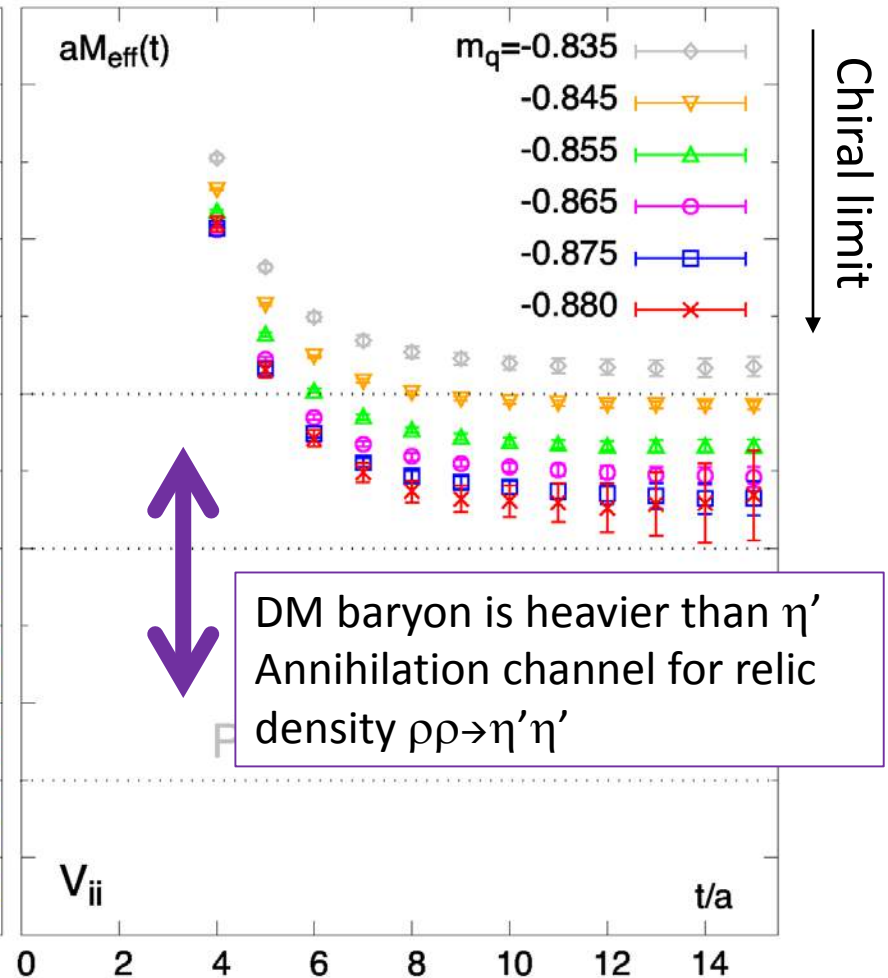
Anthony Francis, R. Jamie Hudspith, Randy Lewis, ST (work in progress)



# Pseudoscalar $\eta'$ mass



# Vector $\rho$ mass (DM)



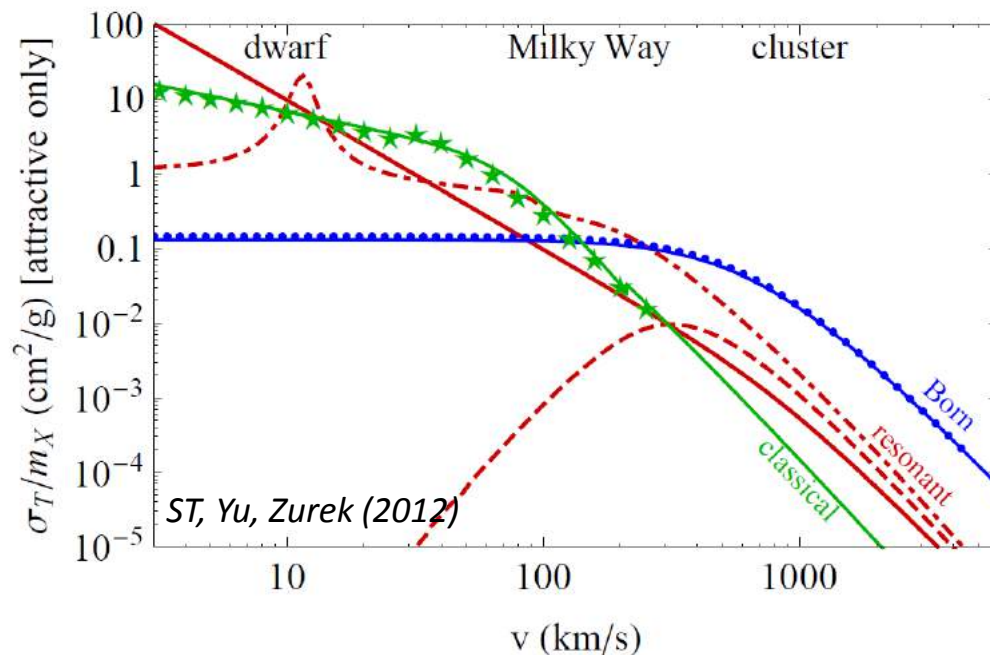
Anthony Francis, R. Jamie Hudspith, Randy Lewis, ST (work in progress)

# Light mediator models

- Weakly-coupled theories
- MSIDM

$$\mathcal{L}_{\text{int}} = \begin{cases} g_\chi \bar{\chi} \gamma^\mu \chi \phi_\mu & \text{(vector mediator)} \\ g_\chi \bar{\chi} \chi \phi & \text{(scalar mediator)} \end{cases}$$

Three parameters:  
 masses  $m_\chi$   $m_\phi$   
 dark fine structure  
 constant  $\alpha_\chi \equiv g_\chi^2/4\pi$



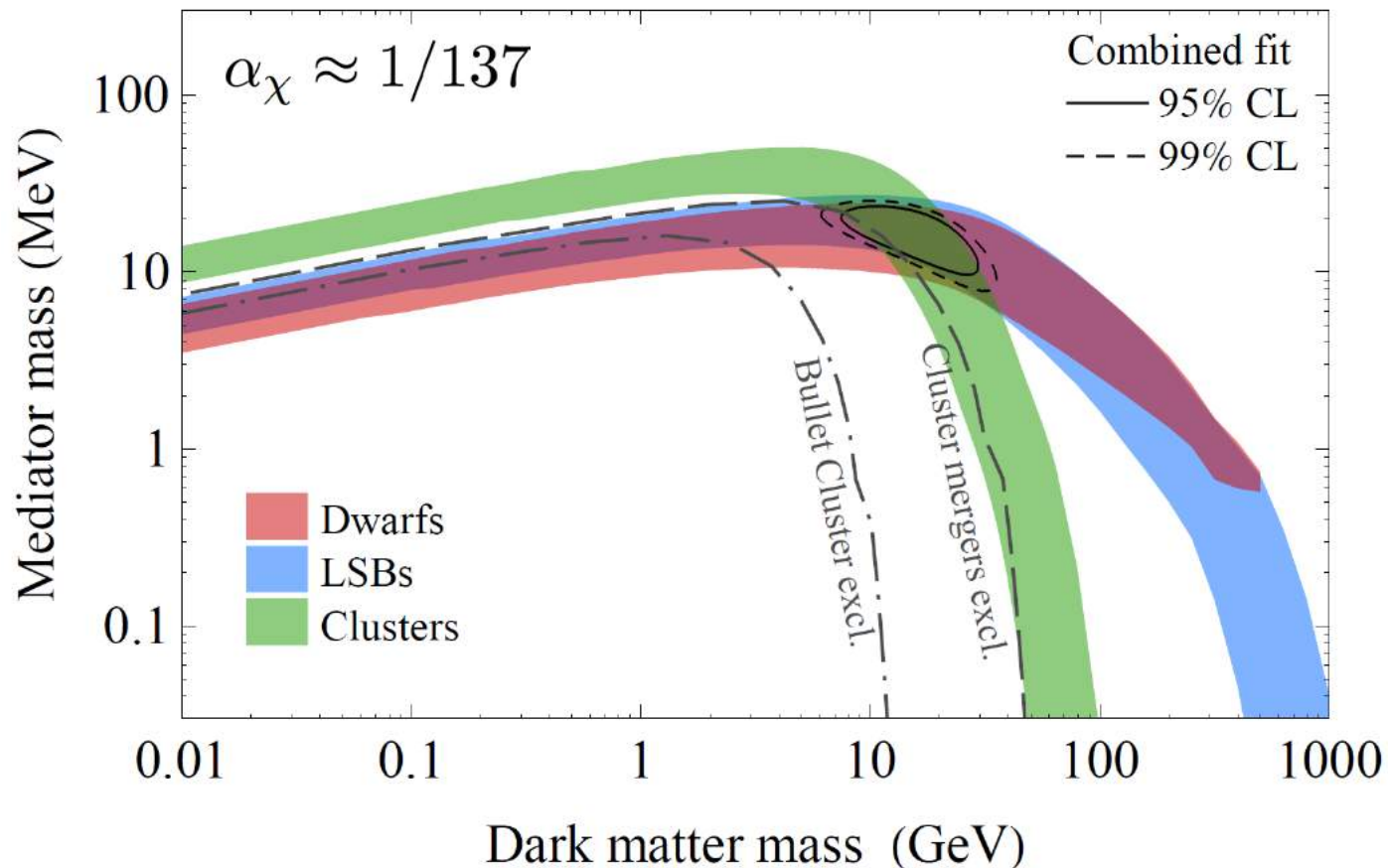
Transfer cross section for  
 different parameter values

Wide range of velocity  
 dependence, but typically  
 suppressed at high velocity

# Light mediator models

Fit parameters from velocity dependence  $\langle\sigma(v)v\rangle/m$  for galaxies and clusters *for a given model*

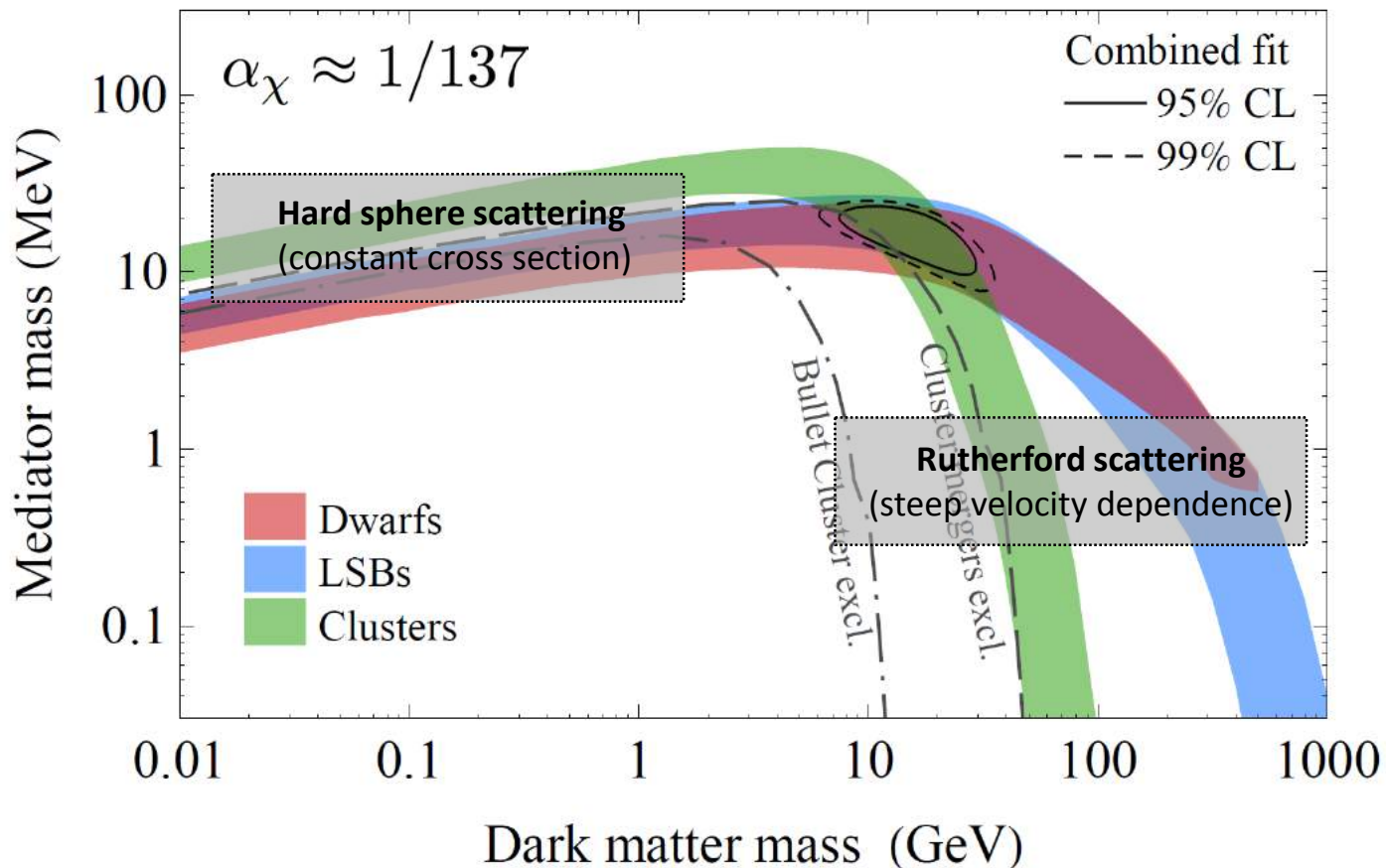
Kaplinghat, ST, Yu (2015)



# Light mediator models

Fit parameters from velocity dependence  $\langle\sigma(v)v\rangle/m$  for galaxies and clusters *for a given model*

Kaplinghat, ST, Yu (2015)



# Portals for light mediators

Mediator may couple to visible sector (SM particles)

Both vector and scalar  $\phi$  may couple via renormalizable interactions. Why not?

Vector mediator case

$$\mathcal{L}_{\text{mixing}} = \frac{\varepsilon_\gamma}{2} \phi_{\mu\nu} F^{\mu\nu}$$

Mixing with photon

Scalar mediator case

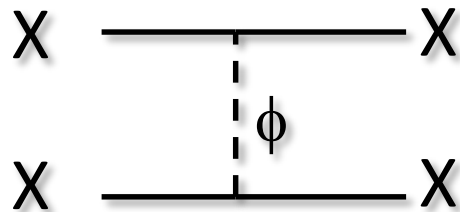
$$V(H, \phi) \supset (a\phi + b\phi^2)|H|^2$$

Mixing with Higgs boson

Cosmological useful

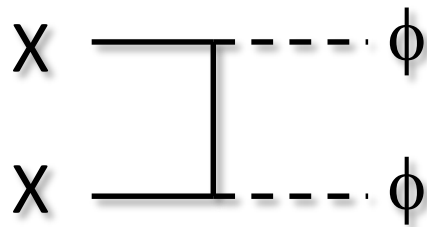
Thermal bath of mediators  $\phi$  can dominate over SIDM component  
Coupling allows mediator decay  $\phi \rightarrow e^+e^-$  to deplete this density

# Self-interacting dark matter paradigm



Self-interactions

DM particle  $X$  + mediator particle  $\phi$



Annihilation

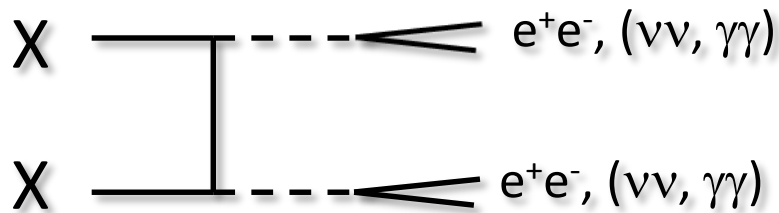
Set relic density  
via freeze-out

$$\phi \xrightarrow{\text{Decay}} \text{SM}$$

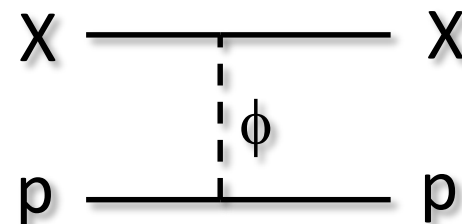
(Deplete  $\phi$  density)

$$\text{SM} \rightarrow \phi \rightarrow \text{SM}$$

Dark photon searches



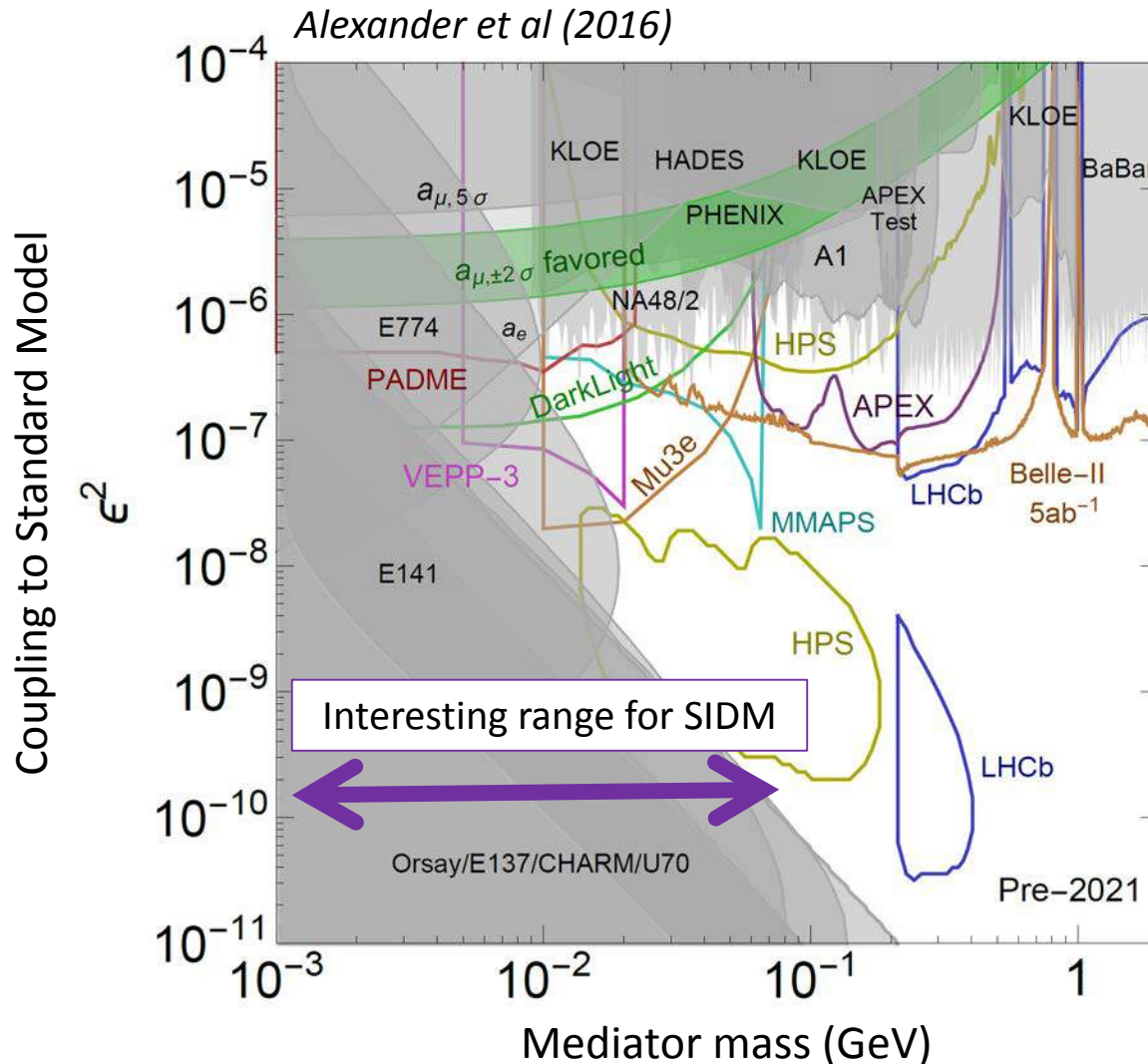
Indirect detection



Direct detection

Capture in sun/earth

# Dark photon searches



Directly produce light mediator  $\phi$  through its coupling to SM particles

e.g.,  $e^-$  beam fixed target experiments

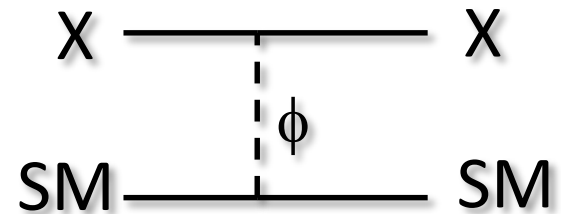
$e^- p \rightarrow e^- p \phi (\rightarrow e^+ e^-)$



# Direct detection

Strongest sensitivity to SIDM coupling  
with Standard Model particles

*Kaplinghat, Tulin, Yu (2013); Del Nobile et al (2015)*



Vector mediator case with kinetic mixing

Spin-independent proton-DM cross section (momentum transfer  $q^2 = 0$ )

$$\sigma_{Xp}^{\text{SI}} \approx 1.5 \times 10^{-24} \text{ cm}^2 \times \varepsilon_\gamma^2 \times \left( \frac{\alpha_X}{10^{-2}} \right) \left( \frac{m_\phi}{30 \text{ MeV}} \right)^{-4}$$

Present limits: for SIDM above few GeV

Coupling to SM must be very tiny (below  $\sim 10^{-10}$ )

Likely SIDM and Standard Model *not* in thermal contact in early Universe  
(during DM freeze-out)

No hope of directly producing mediators in the laboratory



# Direct detection

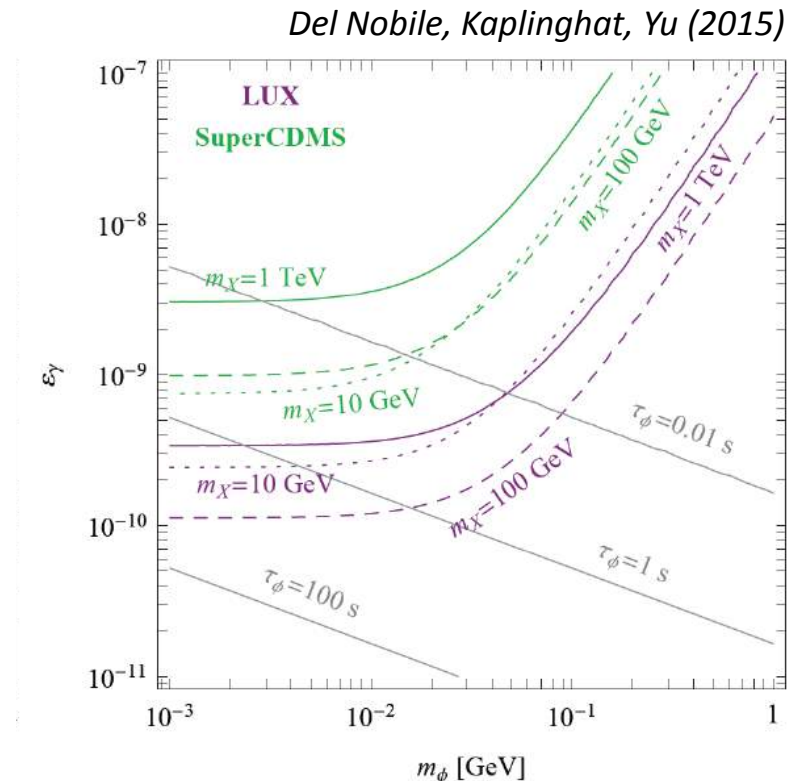
Unique feature of SIDM vs WIMPs:  
Momentum-dependent scattering

Mediator mass  $m_\phi$  can be comparable  
to momentum transfer  $q$

$$\left( \text{Scattering rate} \right) = \left( \text{Scattering rate at } q^2=0 \right) \times \frac{m_\phi^4}{(m_\phi^2 + q^2)^2}$$

Typical momentum transfer for Xenon/Germanium

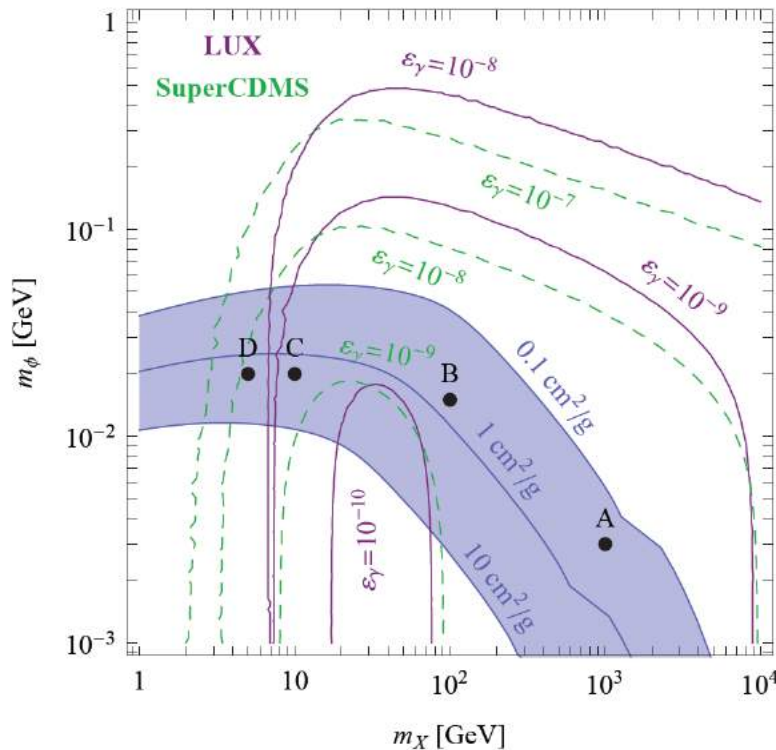
$m_X$ (GeV)	$m_\phi$ (MeV)	$q_{\text{Xe}}$ (MeV)	$q_{\text{Ge}}$ (MeV)
1000	3	127	74
100	15	62	46
10	20	10	10
5	20	5	5



Small improvements in sensitivity can explore large parameter regions for SIDM

# Direct detection limits on SIDM

*Del Nobile et al (2015)*



SIDM above few GeV must couple *very* weakly to the SM

Blue band shows  $\sigma/m$  on dwarf scales preferred by rotation curves and too-big-to-fail

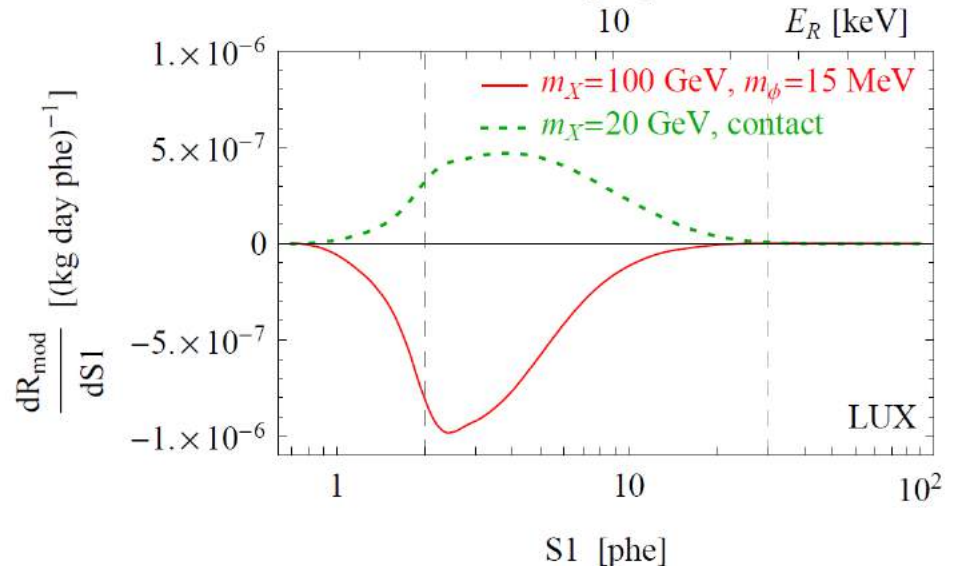
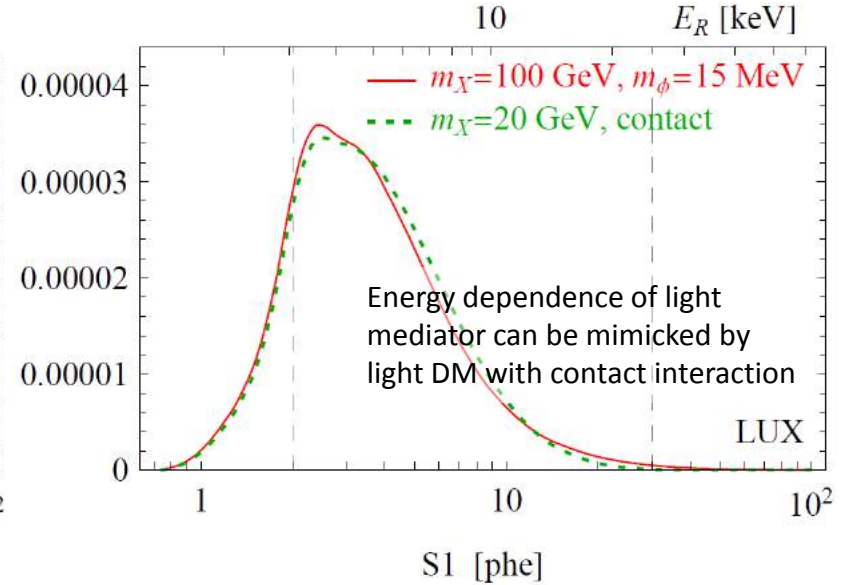
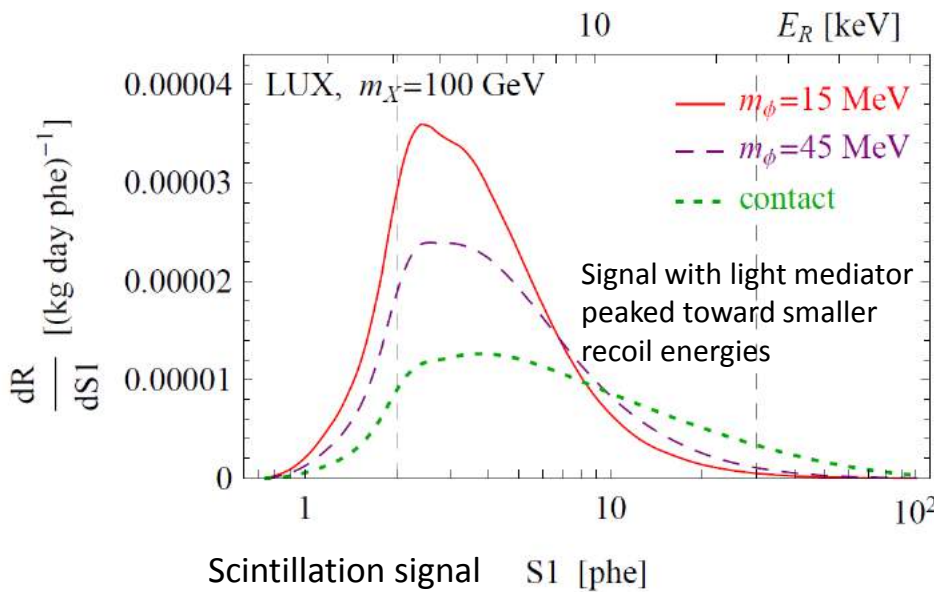
Note: Full LUX exposure is greater by factor  $\sim 3$  (Comparable to Xenon1T)

Lower threshold detectors (SuperCDMS) have greater sensitivity to lighter SIDM mass

*Kahlhoefer, Kulkarni, Wild (2017)*

*See talk by Wild*

# Distinguishing SIDM from WIMPs

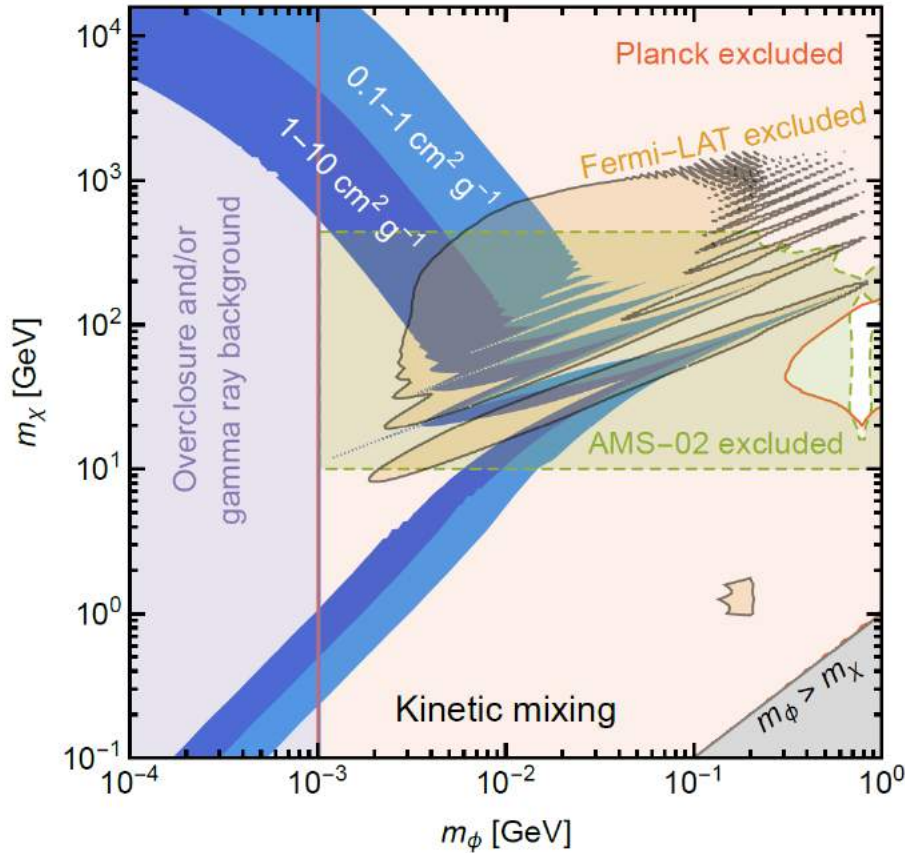


Combination of total rate + annual modulation can distinguish SIDM from WIMPs

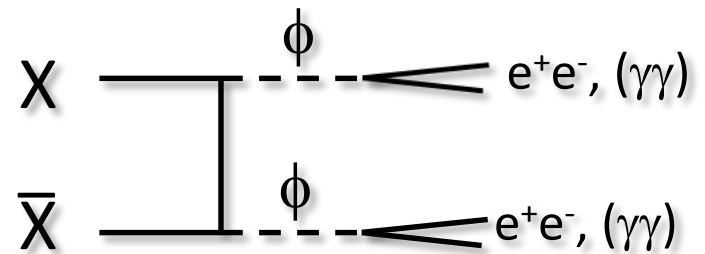
*Del Nobile et al (2015)*

# Indirect detection constraints

*Bringmann, Kahlhoefer, Schmidt-Hoberg, Walia (2016); also Cirelli et al. (2016)*



Visible annihilation signals from SIDM are strongly constrained by CMB



Annihilation for high mass SIDM is boosted by large Sommerfeld factor

Caveats: Asymmetric dark matter, scalar mediator, dark sector decoupled (and cooler) during freeze-out

*See talk by Wild*

# Conclusions

# Dark Matter Self-interactions and Small Scale Structure

Sean Tulin<sup>1,\*</sup> and Hai-Bo Yu<sup>2,†</sup>

Invited review Physics Reports (arXiv:1705.02358)

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### References





$$I(J^{\text{PC}}) = ?(???)$$

Mass unknown

Positive observations	$\sigma/m$	$v_{\text{rel}}$	Observation	Refs.
Cores in spiral galaxies (dwarf/LSB galaxies)	$\gtrsim 1 \text{ cm}^2/\text{g}$	30 – 200 km/s	Rotation curves	[77, 93]
Too-big-to-fail problem				
Milky Way	$\gtrsim 0.6 \text{ cm}^2/\text{g}$	50 km/s	Stellar dispersion	[87]
Local Group	$\gtrsim 0.5 \text{ cm}^2/\text{g}$	50 km/s	Stellar dispersion	[88]
Cores in clusters	$\sim 0.1 \text{ cm}^2/\text{g}$	1500 km/s	Stellar dispersion, lensing	[93, 103]
*** <i>Abell 3827 subhalo merger</i>	$\sim 1.5 \text{ cm}^2/\text{g}$	1500 km/s	DM-galaxy offset	[104]
<i>Abell 520 cluster merger</i>	$\sim 1 \text{ cm}^2/\text{g}$	2000 – 3000 km/s	DM-galaxy offset	[105, 106, 107]
<b>Constraints</b>				
Halo shapes/ellipticity	$\lesssim 1 \text{ cm}^2/\text{g}$	1300 km/s	Cluster lensing surveys	[86]
Substructure mergers	$\lesssim 2 \text{ cm}^2/\text{g}$	$\sim 500 - 4000 \text{ km/s}$	DM-galaxy offset	[92, 108]
Merging clusters	$\lesssim \text{few cm}^2/\text{g}$	2000 – 4000 km/s	Post-merger halo survival (Scattering depth $\tau < 1$ )	Table II
<i>Bullet Cluster</i>	$\lesssim 0.7 \text{ cm}^2/\text{g}$	4000 km/s	Mass-to-light ratio	[81]

\*\*\* We do not use the following data for averages, fits, limits, etc.