## Dark matter cores in galaxies and galaxy clusters

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Workshop "Self-interacting dark matter", NBI, August 3, 2017.



### "Missing mass" problem and DM distribution

Conceptually, dark matter is a solution to the 'missing' (dynamical vs visible) mass problem.

In **outskirts**, visible ('baryonic') mass density falls rapidly with radius, it is (rather) simple to distinguish it from DM.

But in central parts of many objects, baryons dominate:



Courtesy of Seigar, PASP 2008; Newman et al., ApJ 2013b.  $\Rightarrow$  LARGE uncertainty on DM inner slope.

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#### Cores and cusps

Define inner density slope  $\alpha = d\log(\rho_{\rm DM})/d\log(r)$  at  $r \to 0$ 

 $\alpha \sim 1 - \text{cusp}$ , e.g. Navarro-Frenk-White profile

$$\rho_{\rm NFW}(r) = \frac{\rho_s r_s}{r \left(1 + r/r_s\right)^2}$$

(Navarro et al., ApJ 1996, 1997) consistent with DM-only simulations and observations of galaxy clusters.

 $\alpha \sim 0$  means  $\rho_{\text{DM}}(0) = \text{const} - \text{core}$ , e.g. pseudo-isothermal profile

$$\rho_{\rm ISO}(r) = \frac{\rho_c}{1 + r^2/r_c^2}$$

and Burkert profile (Burkert, ApJ 1995)

$$\rho_{\rm BURK}(r) = \frac{\rho_B}{(1 + r/r_B) \left(1 + r^2/r_B^2\right)}$$



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motivated by observations of galaxies.

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#### What can we say regardless of core-cusp problem?

Approximation of cored (ISO, BURK) profiles with cusp (NFW) at intermediate radii (down to  $r_{max}$ ) yields (Boyarsky, Ruchayskiy, DI et al. 2009)

 $r_s \simeq 6.1 r_c; \quad \rho_s \simeq 0.10 \rho_c \qquad (\text{ISO vs NFW})$ 

 $r_s \simeq 1.6 r_B; \quad \rho_s \simeq 0.37 \rho_B \qquad ({\rm BURK \ vs \ NFW})$ 



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#### DM column density

 $S(r_s) \simeq 1.2 \rho_s r_s$  is almost independent from type of DM profile:



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#### Dependence on halo mass

S grows slowly with DM halo mass,  $S \sim (M_{halo})^{\approx 0.2}$  consistent with  $\Lambda$ CDM expectations (Boyarsky, Ruchayskiy, DI et al. 2009)



Also in agreement with the semi-analytical secondary infall model (Boyarsky, Neronov, Ruchayskiy, Tkachev, PRL 2009)



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dSphs are **dispersion-dominated** systems (with no or little rotation). The standard modeling procedure is to solve the **spherically-symmetric** Jeans equations with **unknown** velocity anisotropy (Binney, MNRAS 1980):

$$\frac{d[n(r)\overline{v_r^2}(r)]}{dr} + 2\frac{\beta(r)}{r}n(r)\overline{v_r^2}(r) = -n(r)\frac{GM_{\mathsf{dyn}}(r)}{r^2},$$

n(r) is the stellar number density,  $\overline{v_r^2}(r)$  is the stellar velocity dispersion,  $\beta(r)$  is the stellar velocity anisotropy parameter:

$$\beta(r) = 1 - \frac{\overline{v_{\theta}^2} + \overline{v_{\phi}^2}}{2\overline{v_r^2}}$$

#### Jeans modeling procedure



(Courtesy of El-Badry et al., ApJ 2017)

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#### Velocity anisotropy in dSphs

 $\beta(r)$  usually goes to 0 when  $r \rightarrow 0$ , see e.g. pure-DM dSph (Sparre & Hansen, JCAP 2012), FIRE (EI-Badry et al., ApJ 2017) or APOSTLE simulations (Campbell et al., MNRAS 2017):



 $\beta(r)$  can be **measured** with future proper motion measurements (Gaia – Jin, Helmi & Breddels, 2015, Hubble Astrometry Initiative – Kallivayalil et al., 2015 and Theia – Boehm et al., 2017).

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#### DM mass estimates in dSphs: Jeans equations

Mass inside half-light radius is largely independent on anisotropy, see e.g. Wolf et al., MNRAS 2010: (for 3D half-light radius)



or Walker et al., ApJ 2009 (for 2D half-light radius  $R_{\rm e}$ ):

$$M(R_{\rm e}) = \frac{5\langle\sigma_{\rm los}(R_{\rm e})\rangle^2 R_{\rm e}}{2G_{\rm N}}$$

Image: Image:

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#### DM mass estimates in dSphs: simulations

Recent mass estimates based on APOSTLE simulations (Campbell et al., MNRAS 2017):



and FIRE simulations (Gonzalez-Samaniego et al., 2017).

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#### Multiple stellar populations

If one has two kinematically distinct stellar populations with different  $R_{\rm e}$ , we can use them to measure the slope of M(r), see e.g. Walker & Peñarrubia, ApJ 2011:



For three stellar populations one can derive not only mass (density) slope, but also core radius – e.g. Amorisco et al., MNRAS 2013 found  $r_c = 1.0^{+0.8}_{-0.4}$  kpc in Fornax.

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#### Correlation between 2 subpopulations

APOSTLE simulations yield 2-26% of cusps mimicking as cores (Genina et al., 2017):



Taking this into account decreases inconsistency between Walker & Peñarrubia (2011) measurements and NFW in Sculptor and Fornax to 95.4% (from > 99%) and to 91.8% (from > 96%), respectively.

#### Globular clusters in dSphs

GCs need **cored** DM distribution to survive inside a dSph:



Kleyna et al., ApJ 2003 (Ursa Minor) See also Goerdt et al., MNRAS 2006 (Fornax), Contenta et al., 2017 (Eri II), Amorisco, ApJ 2017 (Eri II & And XXV).



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### Spiral galaxies



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 $V^2(r) = \frac{G_N M(< r)}{r} \Rightarrow \rho(r) = \frac{V_{\rm rot}(r) \left[V_{\rm rot}(r) + 2rV_{\rm rot}'(r)\right]}{G_N r^2} \qquad \qquad \text{Discovery}$ 

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#### Cores in galaxies and galaxy clusters

Radius (arcsec)

## THINGS + LITTLE THINGS (Oh et al. AJ 2015)



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#### What can turn cusp into (apparent) core?

- 'Pressure support' provided by the finite gas velocity dispersion (Pineda et al., MNRAS 2017)
- Non-circular motions in HI disks (Oman et al., 2017)



#### Groups and clusters of galaxies

Dispersion-supported systems, DM is probed by galaxy velocities, strong and weak lensing and X-rays.

Kinematics of galaxy clusters **alone** is consistent with NFW profile.



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#### Galaxy clusters + BGCs

Adding brightest cluster galaxy (BCG) kinematics tends to a cusp flattening to  $\alpha\sim0.5,$  in tension with simulations:



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#### Wobbles in galaxy clusters

Possible solution – BCG can be **offset** of DM halo center. However, for CDM the expected wobbles are much smaller (< 2 kpc) to explain the apparent inconsistency (see David Harvey talk).



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### Summary:

- To look for deviations from LCDM, inner DM slope measuremets are essential;
- For many dark matter-dominated objects, the existence of cores have been reported;
- It results in an open window for alternative (e.g. SIDM) interpretations.

# Thank you for your attention!